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## An improved retrieval of tropospheric nitrogen dioxide from GOME

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[1] We present a retrieval of tropospheric nitrogen dioxide ( $\text{NO}_2$ ) columns from the Global Ozone Monitoring Experiment (GOME) satellite instrument that improves in several ways over previous retrievals, especially in the accounting of Rayleigh and cloud scattering. Slant columns, which are directly fitted without low-pass filtering or spectral smoothing, are corrected for an artificial offset likely induced by spectral structure on the diffuser plate of the GOME instrument. The stratospheric column is determined from  $\text{NO}_2$  columns over the remote Pacific Ocean to minimize contamination from tropospheric  $\text{NO}_2$ . The air mass factor (AMF) used to convert slant columns to vertical columns is calculated from the integral of the relative vertical  $\text{NO}_2$  distribution from a global 3-D model of tropospheric chemistry driven by assimilated meteorological data (Global Earth Observing System (GEOS)-CHEM), weighted by altitude-dependent scattering weights computed with a radiative transfer model (Linearized Discrete Ordinate Radiative Transfer), using local surface albedos determined from GOME observations at  $\text{NO}_2$  wavelengths. The AMF calculation accounts for cloud scattering using cloud fraction, cloud top pressure, and cloud optical thickness from a cloud retrieval algorithm (GOME Cloud Retrieval Algorithm). Over continental regions with high surface emissions, clouds decrease the AMF by 20–30% relative to clear sky. GOME is almost twice as sensitive to tropospheric  $\text{NO}_2$  columns over ocean than over land. Comparison of the retrieved tropospheric  $\text{NO}_2$  columns for July 1996 with GEOS-CHEM values tests both the retrieval and the nitrogen oxide radical ( $\text{NO}_x$ ) emissions inventories used in GEOS-CHEM. Retrieved tropospheric  $\text{NO}_2$  columns over the United States, where  $\text{NO}_x$  emissions are particularly well known, are within 18% of GEOS-CHEM columns and are strongly spatially correlated ( $r = 0.78$ ,  $n = 288$ ,  $p < 0.005$ ). Retrieved columns show more  $\text{NO}_2$  than GEOS-CHEM columns over the Transvaal region of South Africa and industrial regions of the northeast United States and Europe. They are lower over Houston, India, eastern Asia, and the biomass burning region of central Africa, possibly because of biases from absorbing aerosols. **INDEX TERMS:** 0394 Atmospheric Composition and Structure: Instruments and techniques; 0365 Atmospheric Composition and Structure: Troposphere—composition and chemistry; 0345 Atmospheric Composition and Structure: Pollution—urban and regional (0305)

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### 1. Introduction

[2] Nitrogen oxide radicals ( $\text{NO}_x \equiv \text{NO} + \text{NO}_2$ ) originating from combustion, lightning, and soils largely control

tropospheric ozone production [Kasibhatla *et al.*, 1991; Penner *et al.*, 1991; Murphy *et al.*, 1993; Jacob *et al.*, 1996]. Considerable uncertainty exists in the magnitude and distribution of  $\text{NO}_x$  emissions [Emmons *et al.*, 1997; Lee *et al.*, 1997]. Global mapping of nitrogen dioxide ( $\text{NO}_2$ ) atmospheric concentrations from space could provide critical information for constraining  $\text{NO}_x$  emissions and more generally improve our understanding of tropospheric chemistry [National Research Council, 2001].

[3] The Global Ozone Monitoring Experiment (GOME) instrument [Burrows *et al.*, 1993, 1999a; ESA, 1995] on board the European Remote Sensing-2 satellite provides the capability for continuous global monitoring of  $\text{NO}_2$  atmospheric columns through observation of solar backscatter at 0.31 nm spectral resolution between 423 and 451 nm where

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NO<sub>2</sub> has strong absorption features. The satellite was launched in April 1995 into a 98.5° inclination Sun-synchronous orbit, crossing the equator at 1030 AM local time in the descending node. The GOME instrument observes the atmosphere in the nadir view with a surface spatial resolution of 40 km latitude by 320 km longitude in the forward scan, using a scanning mirror to measure 3 such scenes across the flight track. Global coverage is achieved every 3 days after 43 orbits. The instrument has broadband polarization monitoring devices (PMDs) with surface spatial resolution of 20 by 40 km<sup>2</sup> that are useful in determining cloud fraction [Kurosu *et al.*, 1999].

[4] *Leue et al.* [2001] and *Velders et al.* [2001] have previously demonstrated the usefulness of NO<sub>2</sub> measurements from GOME for mapping surface emissions of NO<sub>x</sub>. Over NO<sub>x</sub> source regions the tropospheric component of the NO<sub>2</sub> column is comparable in magnitude to the stratospheric component that originates from N<sub>2</sub>O oxidation in the middle stratosphere. The tropospheric NO<sub>2</sub> column tracks surface NO<sub>x</sub> emissions on a regional scale since NO<sub>2</sub> typically accounts for ~80% of NO<sub>x</sub> in the boundary layer and the lifetime of NO<sub>x</sub> against oxidation in the boundary layer is on the order of a day.

[5] Four major challenges are involved in quantifying tropospheric NO<sub>2</sub> columns from GOME. The first is to determine the total NO<sub>2</sub> slant column from the radiance and solar irradiance measurements. The second is to remove a daily varying artificial offset in the NO<sub>2</sub> slant columns thought to be introduced by the diffuser plate of the GOME instrument. The third is to remove the stratospheric column to obtain the tropospheric residual. The fourth challenge is to convert the tropospheric slant columns to vertical columns with a proper air mass factor (AMF) that accounts for atmospheric scattering, including the effect of clouds [Koelemeijer and Stammes, 1999; Velders *et al.*, 2001; Richter and Burrows, 2002].

[6] Previous retrievals of tropospheric NO<sub>2</sub> [Leue *et al.*, 2001; Velders *et al.*, 2001; Richter and Burrows, 2002] employed differential optical absorption spectroscopy [e.g., Platt, 1994] to determine the NO<sub>2</sub> slant column. Only Richter and Burrows [2002] addressed the diffuser plate artifact by fitting all GOME observations with a single solar spectrum. The separation of tropospheric and stratospheric columns differed slightly among the three studies. Both Leue *et al.* [2001] and Velders *et al.* [2001] used GOME observations of cloudy scenes over marine regions at least 200 km from shore to determine a stratospheric background and infer the global stratospheric distribution using a two-dimensional interpolation algorithm. Richter and Burrows [2002] determined the stratospheric background from GOME observations between 180°W and 170°W and inferred the global stratospheric background assuming zonal invariance. In the radiative transfer calculation used to derive the AMF, all three studies assumed the tropospheric NO<sub>2</sub> column to be confined below about 1.5 km and evenly distributed there. Leue *et al.* [2001] employed land albedos of about 10–20% and ocean albedos of about 5–10% while Velders *et al.* [2001] and Richter and Burrows [2002] assumed surface albedos of 5%. Both Leue *et al.* [2001] and Velders *et al.* [2001] multiplied their tropospheric columns by a correction factor of about 4 over nondesert regions, assuming that clouds obscure the NO<sub>2</sub> column.

Velders *et al.* [2001] estimated an uncertainty of 50% from their assumption of a fixed NO<sub>2</sub> profile in the AMF calculation and an uncertainty of 100% from their cloud correction. Richter and Burrows [2002] examined the differences between scenes with cloud fractions less than 0.1 and greater than 0.3 to infer the amount of NO<sub>2</sub> in the free troposphere.

[7] The present work improves on these retrievals in several aspects. Radiance spectra are directly fitted without high-pass filtering or spectral smoothing following the works of Chance [1998] and Chance *et al.* [2000]. The fit is performed using NO<sub>2</sub> absorption spectra at 293 K, appropriate for boundary layer NO<sub>2</sub>. Correction for spectral undersampling by the GOME instrument, wavelength calibration with a Fraunhofer reference spectrum [Chance and Spurr, 1997], treatment of the Ring effect, and inclusion of a common mode spectrum follow the works of Chance [1998] and Chance *et al.* [2000]. We determine the stratospheric column and the diffuser plate artifact from GOME over Pacific regions with minimal tropospheric NO<sub>2</sub>, correct the latitudinally varying bias introduced by this approach, and estimate the error from the assumption of zonal invariance using limb observations of stratospheric NO<sub>2</sub> from the Halogen Occultation Experiment (HALOE) instrument [Russell *et al.*, 1993]. Our AMF calculation combines a radiative transfer model with local surface albedos determined from GOME and vertical shape factors of NO<sub>2</sub> locally determined from a global 3-D model of tropospheric chemistry, following the formulation of Palmer *et al.* [2001]. We extend that formulation to account for scattering by clouds using local cloud fraction, cloud top pressure, and cloud optical thickness information from the GOME Cloud Retrieval Algorithm (GOMECAT) [Kurosu *et al.*, 1999]. In the GOMECAT algorithm, cloud fraction is determined from the PMDs while cloud top pressure and cloud optical thickness are obtained from GOME radiances in and around the O<sub>2</sub> A band.

[8] We provide below a description of the method and apply it to a retrieval study for July 1996. Section 2 describes the atmospheric chemistry model used in the retrieval. In section 3 we present the AMF formulation to account for scattering by clouds. The fitting of total slant columns is described in section 4. We remove the nontropospheric column in section 5. Tropospheric slant and vertical columns are presented in section 6. Section 7 discusses the errors in the retrieval.

## 2. Atmospheric Chemistry Model Used in the Retrieval

[9] Retrieving tropospheric NO<sub>2</sub> columns from GOME requires some assumptions regarding the vertical distribution of NO<sub>2</sub>. A global 3-D model of tropospheric chemistry is the best source for this information considering the sparseness of NO<sub>2</sub> in situ observations and the large spatial variability of NO<sub>2</sub> profiles. We use the Global Earth Observing System (GEOS)-CHEM model [Bey *et al.*, 2001a] driven by assimilated meteorological observations for 1996, updated every 3–6 hours, from the GEOS of the NASA Data Assimilation Office (DAO) [Schubert *et al.*, 1993]. The model version used here has 26 vertical levels on a sigma coordinate (surface to 0.1 hPa), and a horizontal

**Table 1.** Annual Global NO<sub>x</sub> Emissions in the GEOS-CHEM Model

Source	Emission Rate, Tg N yr <sup>-1</sup>
Fossil fuel combustion	23.1
Soils	5.2
Biomass burning	5.1
Lightning	2.9
Biofuels	2.2
Aircraft	0.5
Stratosphere	0.2 <sup>a</sup>

<sup>a</sup>The cross-tropopause NO<sub>y</sub> flux is 0.7 Tg N yr<sup>-1</sup> (including 0.2 Tg N yr<sup>-1</sup> as NO<sub>x</sub> and 0.5 Tg N yr<sup>-1</sup> as HNO<sub>3</sub>).

resolution of 2° latitude by 2.5° longitude. It includes a detailed description of tropospheric ozone–NO<sub>x</sub>–hydrocarbon chemistry. It solves the chemical evolution of about 120 species with a Gear solver [Jacobson and Turco, 1994] and transports 24 tracers. Photolysis rates are computed using the Fast-J radiative transfer algorithm [Wild *et al.*, 2000] which includes Rayleigh scattering as well as Mie scattering by aerosols and clouds. The annual mean tropopause is diagnosed locally in the model using the standard criterion of a 2 K km<sup>-1</sup> lapse rate. This model version (based on GEOS-CHEM 4.11, <http://www-as.harvard.edu/chemistry/trop/geos>) includes several updates relative to the original Bey *et al.* [2001a] version, as described by Martin *et al.* [2002]. The most important for the present application are monthly averaged UV surface reflectivity fields [Herman *et al.*, 1997], Mie scattering by mineral dust, heterogeneous chemistry on mineral dust aerosols, improved biomass burning and biofuel emission inventories, and improved seasonal variation in biomass burning emissions as summarized below.

[10] Table 1 contains the annual global NO<sub>x</sub> emissions used in the model. Anthropogenic NO<sub>x</sub> emissions are from the Global Emission Inventory Activity (GEIA) [Benkovitz *et al.*, 1996] partitioned among individual countries and scaled to 1996 levels as described by Bey *et al.* [2001a]. Emissions of NO<sub>x</sub> from lightning are linked to deep convection following the parameterization of Price and Rind [1992] as implemented by Wang *et al.* [1998] with vertical profiles from the work of Pickering *et al.* [1998]. Soil NO<sub>x</sub> emissions are computed locally using a modified version of the Yienger and Levy [1995] algorithm, as described by Wang *et al.* [1998] and Bey *et al.* [2001a]. New emission inventories are used for biofuels and biomass burning (J. Logan and R. Yevich, personal communication, 2001). We use vegetation specific emission factors as described by Staudt *et al.* (Sources and chemistry of nitrogen oxides over the tropical Pacific, submitted to *Journal of Geophysical Research*, 2002). Seasonal variation in biomass burning emissions is determined from satellite observations [Duncan *et al.*, 2002]. Ship emissions of NO<sub>x</sub> are from GEIA (0.2 Tg N yr<sup>-1</sup>). Corbett *et al.* [1999] have proposed that ship emissions may be much higher (3.0 Tg N yr<sup>-1</sup>), but this would result in a large model overestimate of NO<sub>x</sub> over the North Atlantic [Kasibhatla *et al.*, 2000; Davis *et al.*, 2001].

[11] Of particular interest here is the ability of the model to provide a realistic simulation of the tropospheric NO<sub>2</sub>

relative vertical distribution (shape factor) for the AMF calculation. Few in situ observations of NO<sub>2</sub> exist, but a large body of aircraft observations for NO is available [Emmons *et al.*, 1997; Thakur *et al.*, 1999]. Considering that the NO<sub>2</sub>/NO ratio from photochemical steady state in the model is known to match observations closely [Bradshaw *et al.*, 1999], observed vertical profiles of NO provide a good surrogate for NO<sub>2</sub> evaluation. Detailed evaluations of the GEOS-CHEM NO fields with observations are presented in several papers [Bey *et al.*, 2001a, 2001b; Fiore *et al.*, 2002; Martin *et al.*, 2002]. They show that the model generally captures the spatial and temporal variability in NO profiles, reproducing observed NO concentrations within a factor of 2. The largest relative discrepancies are over the Pacific where the model is up to a factor of 2 too high [Bey *et al.*, 2001a].

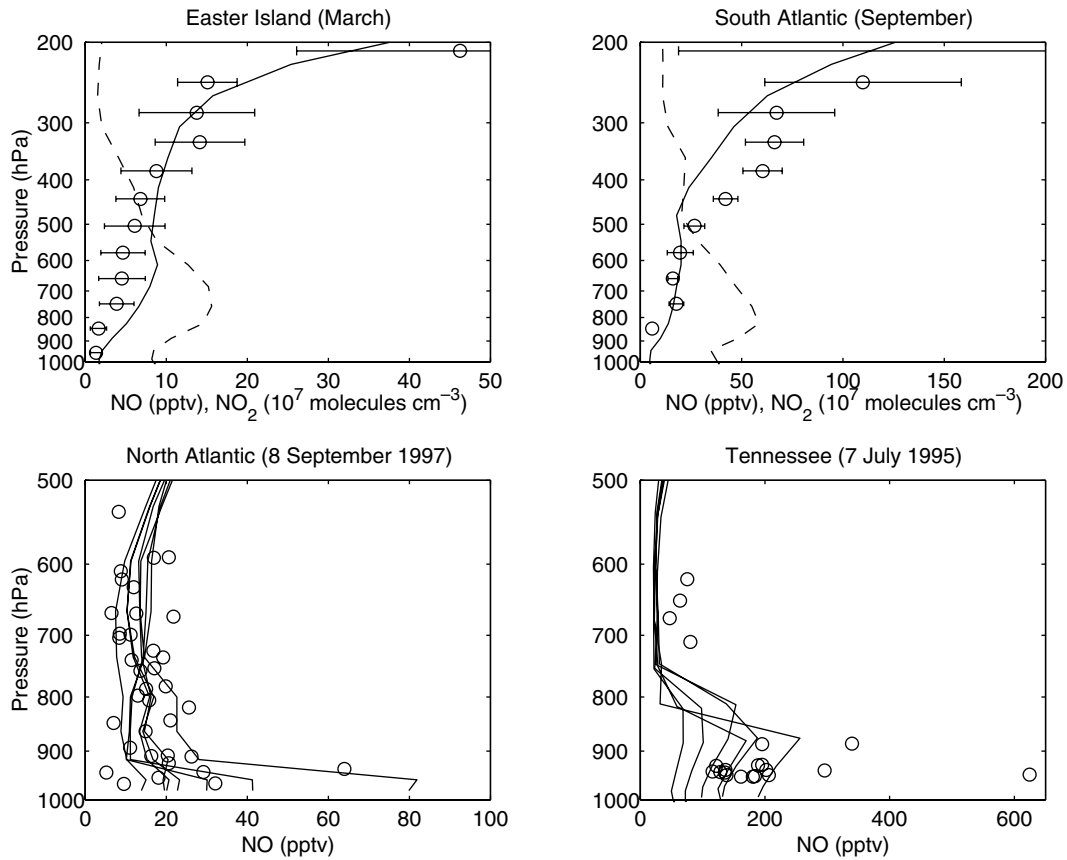
[12] Figure 1 shows four representative profiles over a range of conditions: (top left) clean troposphere, (top right) remote troposphere affected by biomass burning outflow, (bottom left) ocean region downwind of a major source region, and (bottom right) source region during summer. The corresponding NO<sub>2</sub> mixing ratio profiles would be skewed toward the lower troposphere since the daytime NO<sub>2</sub>/NO ratio typically decreases by a factor of 25 from the surface to the upper troposphere [Bradshaw *et al.*, 1999], largely due to the temperature dependence of the NO + O<sub>3</sub> reaction. As shown in Figure 1, the NO<sub>2</sub> number density profiles are skewed further toward the lower troposphere. As a result, GOME is relatively insensitive to upper tropospheric NO<sub>x</sub> enhancements from aircraft and lightning.

[13] The shape of the vertical profile shows large variability depending on the region, and the model largely captures this variability, as shown in Figure 1. The model underestimates the high upper tropospheric NO concentrations over the tropical South Atlantic from biomass burning and lightning [Jacob *et al.*, 1996; Pickering *et al.*, 1996] by up to 50%. Over the Pacific, the model overestimates lower tropospheric NO concentrations by up to 50%. Over the North Atlantic, the model reproduces well the observed vertical profile and high spatial variability. Over Tennessee the model simulates the observed boundary layer enhancement but underestimates its magnitude. The profile shape over source regions is largely determined by the local boundary layer depth. The model simulation of boundary layer depths will be discussed in section 7.

### 3. AMF Calculation

[14] The AMF is defined here as the ratio of the fitted (“slant”) to the vertical tropospheric column of NO<sub>2</sub>. The stratospheric component of the slant column is subtracted prior to the application of the AMF as described in section 5. The AMF is sensitive to the relative vertical distribution of NO<sub>2</sub> due to Rayleigh scattering and to Mie scattering by clouds. Backscattered radiances measured by GOME can be strongly influenced by clouds, even if clouds constitute only a small fraction of a GOME scene [Koelemeijer and Stammes, 1999; Velders *et al.*, 2001; Richter and Burrows, 2002]. An important feature of our AMF formulation is that it enables quantitative retrieval for partly cloudy scenes, which represent the general case for GOME because of the





**Figure 1.** Comparison of NO concentrations between the GEOS-CHEM model (solid lines) and aircraft observations (circles) over Easter Island in March 1999 [Raper *et al.*, 2002], the tropical South Atlantic in September 1992 [Fishman *et al.*, 1996b], the North Atlantic off the east coast of Nova Scotia on 8 September 1997 [Ryerson *et al.*, 1999], and Tennessee on 7 July 1995 [Hübner *et al.*, 1998]. For Easter Island and the tropical South Atlantic, the observations are means and standard deviations from several flights in the region [Bey *et al.*, 2001a]. Model results are monthly means for 1997 (top left) and 1996 (top right), i.e., not the same years as the observations. For the North Atlantic and Tennessee, the model results are for the specific day of the flight and the different lines represent the ensemble of grid squares sampled by the flight tracks [Fiore *et al.*, 2002]. The dashed lines in the top panels show the modeled NO<sub>2</sub> number density profiles.

large scene size ( $40 \times 320 \text{ km}^2$ ). The method is described here for NO<sub>2</sub> and for a single cloud layer in each scene, but can be applied to any optically thin absorber and to multiple cloud or aerosol layers. A scattering aerosol layer in an otherwise clear sky would be diagnosed as a cloud of small optical thickness in the GOMECAT algorithm.

### 3.1. General AMF Formulation

[15] We apply the general AMF formulation of Palmer *et al.* [2001] to tropospheric NO<sub>2</sub> from GOME. This formulation decouples the vertical dependence of the GOME sensitivity to the atmospheric species of interest (calculated with a radiative transfer model) from the shape of the vertical profile of concentrations (calculated with an atmospheric chemistry model). Dimensionless shape factors  $S(\sigma)$  over the sigma ( $\sigma$ ) vertical coordinate are determined from the GEOS-CHEM model for each individual observation scene

$$S(\sigma) = C(\sigma) \frac{\Omega_{\text{air}}}{\Omega} \quad (1)$$

where  $\Omega_{\text{air}}$  and  $\Omega$  are the tropospheric vertical columns of air and NO<sub>2</sub>, and  $C(\sigma)$  is the NO<sub>2</sub> mixing ratio. Pressure  $P$  is related to  $\sigma$  by  $P = \sigma(P_s - P_u) + P_u$ , where  $P_u$  and  $P_s$  respectively represent the pressures at the model upper boundary and at the surface. Scattering weights  $w(\sigma)$  describe the sensitivity of the backscattered radiance  $I$  observed by GOME to the abundance of NO<sub>2</sub> at each  $\sigma$ -level

$$w(\sigma) = -\frac{1}{\text{AMF}_G} \frac{\alpha(\sigma)}{\alpha_e} \frac{\partial(\ln I)}{\partial \tau} \quad (2)$$

where  $\alpha(\sigma)$  is the temperature-dependent absorption cross section ( $\text{m}^2 \text{ molecules}^{-1}$ ),  $\alpha_e$  is the effective absorption cross section ( $\text{m}^2 \text{ molecules}^{-1}$ ) representing a weighted average over the tropospheric column [Palmer *et al.*, 2001], and  $\partial \tau$  is the incremental NO<sub>2</sub> optical depth as a function of  $\sigma$ . The geometric air mass factor  $\text{AMF}_G$ , determined simply from the geometric path correction, normalizes the scattering weight such that  $w(\sigma) = 1$  in a nonscattering

atmosphere. It is a function of the solar zenith angle  $\theta_o$  and satellite viewing angle  $\theta$

$$\text{AMF}_G = \sec\theta_o + \sec\theta \quad (3)$$

The AMF is then given by [Palmer *et al.*, 2001]

$$\text{AMF} = \text{AMF}_G \int_{\sigma_T}^1 w(\sigma) S(\sigma) d\sigma \quad (4)$$

where the integral is taken here from the model tropopause  $\sigma_T$  to the surface. In the absence of scattering, the AMF reduces to  $\text{AMF}_G$ .

### 3.2. Treatment of Partly Cloudy Scenes

[16] Although the AMF formulation described in equations (1)–(4) is applicable to any scattering atmosphere, Palmer *et al.* [2001] calculated the scattering weights,  $w(\sigma)$ , solely for a clear-sky Rayleigh scattering atmosphere. We extend here the AMF formulation to partly cloudy scenes, as typically observed by GOME. The GOMECAT algorithm provides cloud fraction, cloud top pressure, and cloud optical thickness for each scene. The measurement of cloud optical thickness eliminates the need to include a “ghost column” (an assumed value for the column below the cloud) used in other retrievals [i.e., McPeters *et al.*, 1998]. Instead we use a radiative transfer model that includes Mie scattering by clouds to calculate scattering weights for both the clear-sky ( $w_a$ ) and cloudy ( $w_c$ ) fractions of the scene at all levels in the troposphere:

$$w_a(\sigma) = -\frac{1}{\text{AMF}_G} \frac{\alpha(\sigma)}{\alpha_e} \frac{\partial(\ln I_a)}{\partial\tau} \quad (5a)$$

$$w_c(\sigma) = -\frac{1}{\text{AMF}_G} \frac{\alpha(\sigma)}{\alpha_e} \frac{\partial(\ln I_c)}{\partial\tau} \quad (5b)$$

We have decomposed the backscattered radiance  $I$  ( $\text{W m}^{-2} \text{ nm}^{-1} \text{ sr}^{-1}$ ) observed by GOME for the entire scene into the contributions from the clear-sky ( $I_a$ ) and cloudy fractions ( $I_c$ )

$$I = I_a(1-f) + I_c f \quad (6)$$

where  $f$  is the cloud fraction ( $0 \leq f \leq 1$ ). In this manner, AMFs can be calculated for the clear-sky and cloudy fractions of the scene (assuming the same shape factor for each fraction)

$$\text{AMF}_a = \text{AMF}_G \int_{\sigma_T}^1 w_a(\sigma) S(\sigma) d\sigma \quad (7a)$$

$$\text{AMF}_c = \text{AMF}_G \int_{\sigma_T}^1 w_c(\sigma) S(\sigma) d\sigma \quad (7b)$$

[17] The ratio  $I_a/I_c$  can be expressed in terms of the reflectivity of each subscene  $R_a$  and  $R_c$

$$\frac{I_a}{I_c} = \frac{R_a}{R_c} \quad (8)$$

where the reflectivity is defined as

$$R_a = \frac{\pi I_a}{E_o \cos\theta_o} \quad (9a)$$

$$R_c = \frac{\pi I_c}{E_o \cos\theta_o} \quad (9b)$$

and  $E_o$  is the solar irradiance at the top of the atmosphere perpendicular to the direction of incident sunlight. The reflectivity of each subscene includes contributions from surface albedo, Rayleigh scattering, and also cloud scattering for the cloudy subscene. Values of  $R_a$  and  $R_c$  can be obtained from a radiative transfer model as described below. Substitution of equations (6) and (8) into equation (2) shows that the actual scattering weights  $w(\sigma)$  for the partly cloudy scene are the averages of  $w_a(\sigma)$  and  $w_c(\sigma)$  weighted by the contribution of each subscene to the backscattered radiance observed by the satellite

$$w(\sigma) = \frac{w_a(\sigma) R_a (1-f) + w_c(\sigma) R_c f}{R_a (1-f) + R_c f} \quad (10)$$

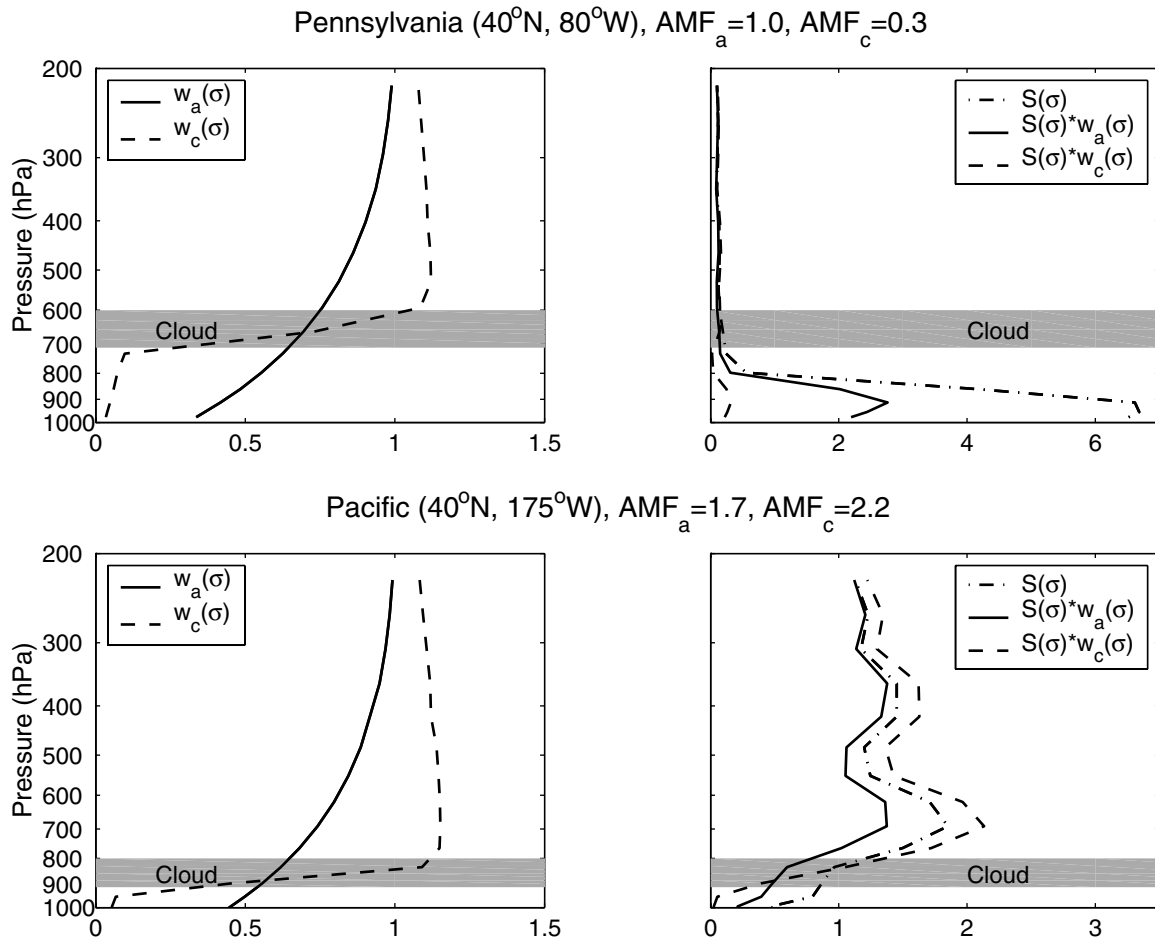
Substituting equation (10) into equation (4) yields a similar relationship for the AMF

$$\text{AMF} = \frac{\text{AMF}_a R_a (1-f) + \text{AMF}_c R_c f}{R_a (1-f) + R_c f} \quad (11)$$

### 3.3. Application to Retrievals of NO<sub>2</sub> From GOME

[18] We calculate scattering weights ( $w_a$ ,  $w_c$ ) and subscene reflectivities ( $R_a$ ,  $R_c$ ) using the Linearized Discrete Ordinate Radiative Transfer (LIDORT) model [Spurr *et al.*, 2001]. The LIDORT model solves the radiative transfer equation in a multilayer atmosphere with multiple scattering using the discrete ordinate method [Chandrasekhar, 1960]. We calculate the scattering weights on an altitude coordinate with 0.5 km vertical resolution below 18 km and coarser resolution from 18 km to the top of the atmosphere (65 km). We then map the results onto the GEOS-CHEM  $\sigma$ -coordinate. Rayleigh scattering cross sections are calculated as in the work of Chance and Spurr [1997]. Vertical profiles of temperature, pressure, and ozone used in LIDORT are from a midlatitude summer *US Standard Atmosphere* [1976]. Using profiles from tropical or winter atmospheres affects the AMFs by less than 0.2%.

[19] The surface albedo is treated as Lambertian. We use a surface albedo database derived from GOME measurements at 440 nm [Koelmeijer *et al.*, 2002], generated as follows. Effective scene albedo was determined for each GOME measurement for the month of July for the years 1995–2000, using the Doubling-Adding KNMI radiative transfer code [De Haan *et al.*, 1987; Stammes, 2000]. The effective scene albedo is the calculated Lambertian surface albedo required to match the observed reflectivity at the top of the atmosphere, assuming a Rayleigh scattering atmosphere. The effective scene albedos were binned by month



**Figure 2.** Derivation of the AMFs for sample GOME viewing scenes over western Pennsylvania (40°N, 80°W) and the central North Pacific (40°N, 175°W). The left panels show scattering weights for the clear-sky  $w_a$  and cloudy  $w_c$  fractions of the scenes as determined with the LIDORT radiative transfer model. The cloud optical thickness assumed for the cloudy fractions is 10 in both scenes. The cloud top pressure is 800 hPa for the Pacific scene and 600 hPa for the Pennsylvania scene. The right panels show the shape factors  $S(\sigma)$  from the GEOS-CHEM model and the products of the shape factor and scattering weight whose integration by equation (7) yields the AMFs. In both scenes, AMF<sub>G</sub> = 2.1.

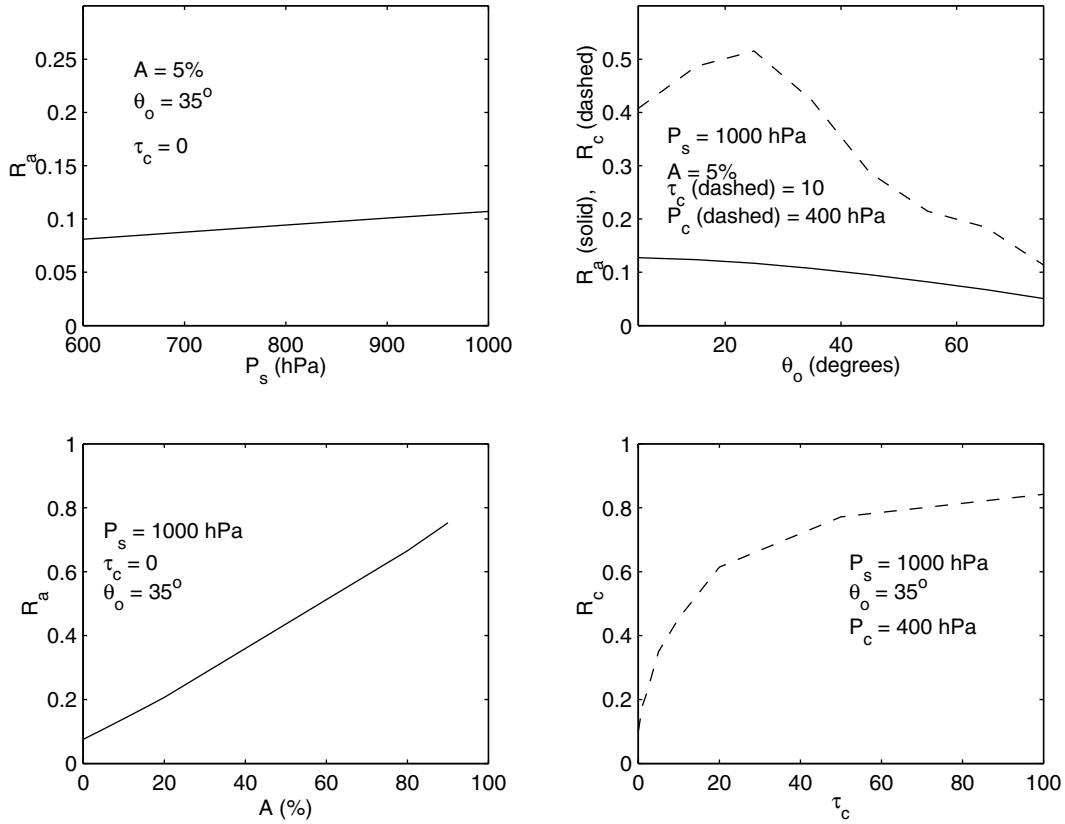
and in grid-cells of 1° by 1°. The surface albedo was then determined as the minimum effective scene albedo in each grid-cell and each month. Effects of persistent clouds over ocean were reduced by replacing the values in such grid-cells with a weighted average of adjacent grid-cells.

[20] The surface albedo is generally between 2% and 10%, except over deserts, snow and ice, and regions of persistent clouds that bias the measurement of surface albedo. Cloud top pressure and cloud optical thickness are obtained from GOMECAT. We estimate cloud base by assuming a cloud optical thickness increment of 8 for each 100 hPa of cloud [Hansen *et al.*, 1983]. The cloud droplet scattering phase function is based on Mie calculations as provided in Fast-J [Wild *et al.*, 2000] for a gamma distribution of cloud droplet sizes with mode radius of 8 μm. The single scattering albedo of cloud droplets is assumed to be 1.

[21] Both  $R_a$  and  $w_a$  are dependent on wavelength  $\lambda$ , surface pressure  $P_s$ , surface albedo  $A$ , solar zenith angle  $\theta_o$ , and viewing zenith angle  $\theta$ . Additionally  $w_c$  and  $R_c$  are

dependent on cloud fraction  $f$ , cloud optical thickness  $\tau_c$ , and cloud top pressure  $P_c$ . We tabulate these dependencies in the center of the NO<sub>2</sub> fitting window (437 nm) for different  $\theta_o$  (5°, 15°, 25°, 35°, 45°, 55°, 65°, 75°, 85°),  $\theta$  (0°, 23°),  $A$  (0, 0.05, 0.10, 0.15, 0.20, 0.80, 0.90),  $P_s$  (1000, 900, 800, 600 hPa),  $P_c$  (900, 800, 600, 400, 200 hPa), and  $\tau_c$  (0, 1, 2, 5, 10, 20, 50, 100). Calculations are averaged over azimuth angle. Wavelength and azimuth angle dependences are negligible over the range of the fitting window (423–451 nm) and GOME viewing angles (0°–23°).

[22] Figure 2 shows  $w_a(\sigma)$ ,  $w_c(\sigma)$ ,  $S(\sigma)$ , and their products for sample scenes in the nadir view ( $\theta = 0^\circ$ ) over the central North Pacific and western Pennsylvania. We see that  $w_a(\sigma)$  is slightly lower over land than over ocean because of the lower surface albedo of vegetated land surfaces around 440 nm. The scattering weights can be regarded as sensitivities relative to a nonscattering atmosphere. For example, Figure 2 for the clear-sky scene over the Pacific shows that GOME has a sensitivity of roughly 45% to the mixing ratio



**Figure 3.** Solar reflectivity at 437 nm for a nadir viewing satellite instrument for clear sky ( $R_a$ ) and cloudy sky ( $R_c$ ) (equations (9)). The plots show the dependencies of  $R_a$  and  $R_c$  on surface pressure  $P_s$ , solar zenith angle  $\theta_o$ , surface albedo  $A$ , and cloud optical thickness  $\tau_c$ , with all other variables referenced to a standard case ( $P_s = 1000$  hPa,  $\theta_o = 35^\circ$ ,  $A = 5\%$ ,  $\tau_c = 0$  or 10, and cloud top pressure  $P_c = 400$  hPa).

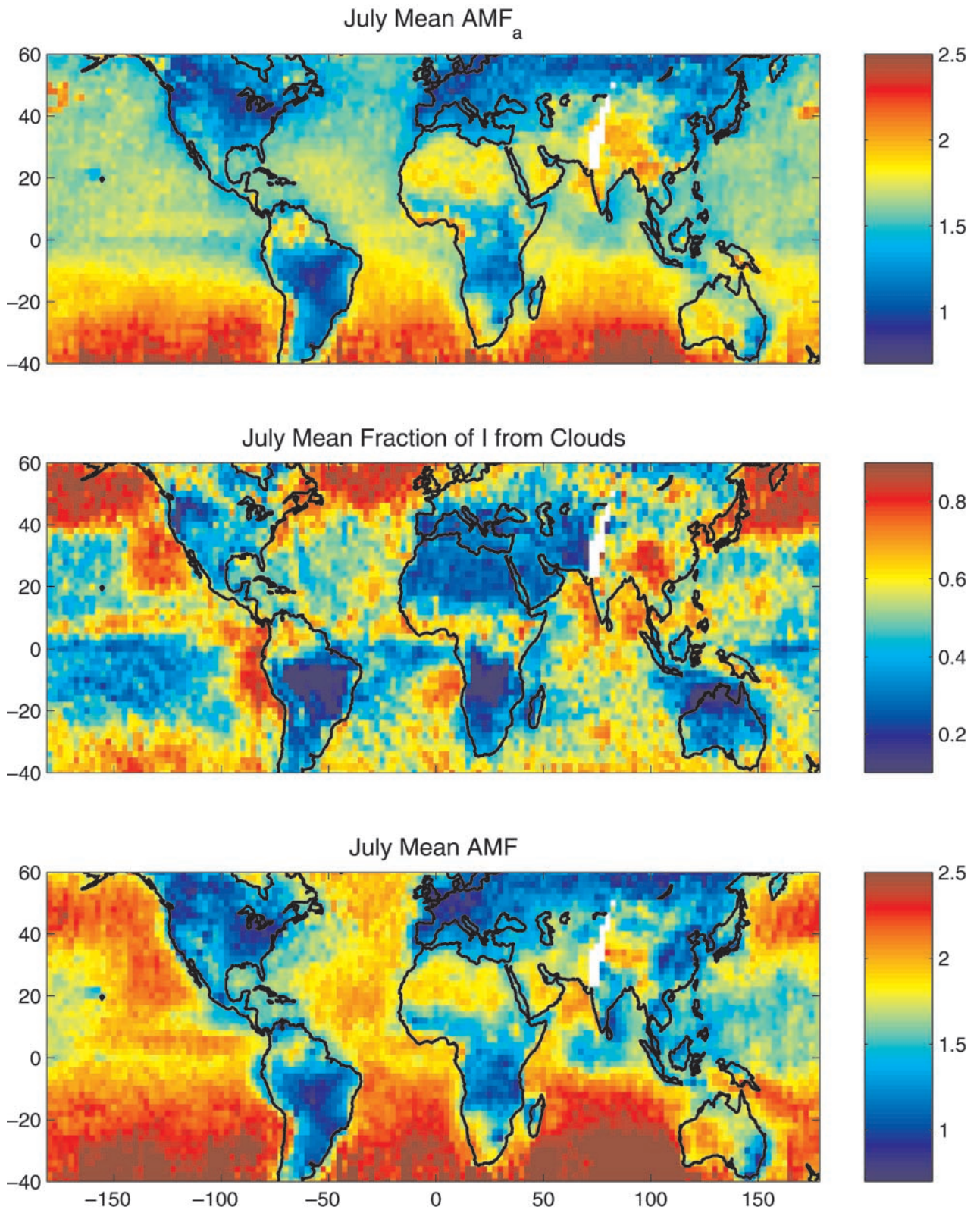
of NO<sub>2</sub> near the surface as compared to 65% at 800 hPa, and  $\sim 100\%$  at 300 hPa. Calculations of  $w_c(\sigma)$  are shown for illustrative cloud top pressures of 800 hPa for the Pacific scene (stratus) and 600 hPa for the Pennsylvania scene. Both clouds have an optical thickness of 10 and hence a pressure thickness of 125 hPa. Above the cloud  $w_c(\sigma)$  is enhanced with respect to  $w_a(\sigma)$  but the opposite is true below the cloud [Koelemeijer and Stammes, 1999; Richter and Burrows, 2002]. Above the cloud  $w_c(\sigma)$  is larger than 1 due to multiple scattering. Little sensitivity to cloud top pressure is exhibited by  $w_c(\sigma)$  above cloud top.

[23] The shape factors in Figure 2 illustrate the range of vertical distributions of NO<sub>2</sub>: minima in the boundary layer over remote oceans (bottom right) and maxima in the boundary layer over continental source regions (top right). The product  $w_a(\sigma)S(\sigma)$  indicates that slant columns observed by GOME in clear-sky conditions are mostly determined by NO<sub>2</sub> in the free troposphere over remote oceans, and by NO<sub>2</sub> in the boundary layer over continental source regions. For clear-sky conditions in both cases,  $AMF_a$  is less than  $AMF_G$  because of the decrease in sensitivity toward the surface. For cloudy conditions, the slant columns observed by GOME are primarily from NO<sub>2</sub> within and above the cloud. As a result  $AMF_c$  can be larger than  $AMF_a$  if little NO<sub>2</sub> is below the cloud (bottom right) or can be smaller than  $AMF_a$  if the cloud obscures boundary layer NO<sub>2</sub> (top right).

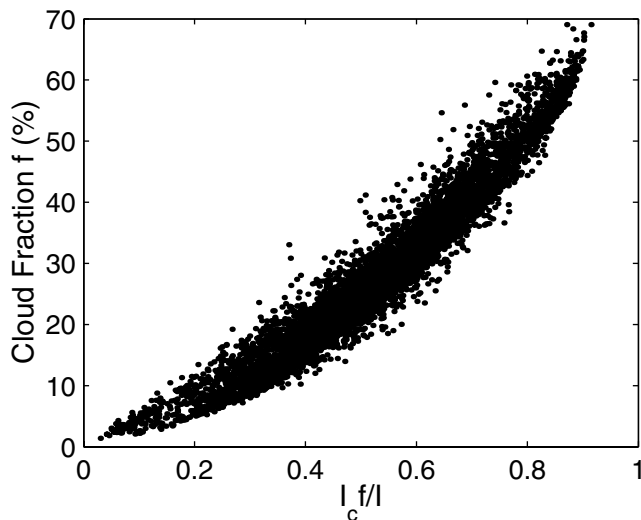
[24] Figure 3 shows the dependence of the clear-sky and cloudy reflectivities ( $R_a$  and  $R_c$ ) on surface pressure  $P_s$ , solar zenith angle  $\theta_o$ , surface albedo  $A$ , and cloud optical thickness  $\tau_c$ , again for an illustrative case. For low surface albedo  $R_a$  increases linearly with surface pressure due to increasing Rayleigh scattering optical depth. The dependence of  $R_a$  on surface pressure decreases with increasing surface albedo, becoming nearly independent of surface pressure for a surface albedo greater than 80% (not shown). Similarly,  $R_c$  exhibits little variation with cloud top pressure (not shown). The top right panel in Figure 3 shows the decrease of  $R_a$  with increasing  $\theta_o$ . Although Rayleigh optical depth increases with  $\theta_o$ , the Rayleigh scattering phase function is forward and backward peaked, scattering less radiation sideways toward the nadir when  $\theta_o$  increases. The phase function for Mie scattering cloud droplets is more strongly peaked with more angular features, hence the more pronounced relationship between  $R_c$  and  $\theta_o$ . The bottom left panel in Figure 3 shows that  $R_a$  has a near-linear relationship with  $A$ . Calculations for a zero surface albedo show that Rayleigh scattering alone produces  $R_a$  of about 10%. Figure 3 (bottom right) illustrates that  $R_c$  is a strong function of  $\tau_c$ , especially for  $\tau_c$  less than 20, beyond which  $R_c$  approaches an asymptotic value. The asymptotic value decreases as  $\theta_o$  increases.

[25] We calculate the AMF for each individual GOME observation scene in July 1996 as a function of local  $P_s$ ,  $A$ ,  $\theta_o$ ,  $\theta_f$ ,  $\tau_c$ , and  $P_c$ , using tabulated values of ( $w_a$ ,  $w_c$ ) and





**Figure 4.** (top) Clear-sky AMF (equation (7a)). (middle) Fraction of GOME backscattered radiance contributed by clouds (equation (12)). (bottom) Actual AMF accounting for clouds (equation (11)). Values are means for July 1996. There are no GOME data for the white area over central Asia due to the absence of GOME observations during normal ERS-2 tape recorder operations.



**Figure 5.** Relationship between GOMECAT cloud fraction and the fraction of the backscattered radiance from clouds (equation (12)) for the ensemble of GOME scenes in July 1996.

( $R_a$ ,  $R_c$ ) as described above together with equations (7) and (11). We discard scenes with an AMF less than 0.5 (less than 0.1% of the total observations) to remove outliers with high cloud fraction, low cloud top pressure, and high optical thickness that obscure the bulk of tropospheric NO<sub>2</sub>. Global AMF<sub>a</sub> values (Figure 4, top) are generally in the range 0.8–1.5 over vegetated land surfaces, and 1.5–2.5 over ocean. Surface NO<sub>x</sub> emissions from anthropogenic activity, soils, and biomass burning introduce a steep vertical gradient in NO<sub>2</sub> concentrations almost everywhere over land; NO<sub>2</sub> profiles over oceans typically show less vertical gradient (Figure 2). Arid (deep boundary layers and high surface albedos) and intensely convective continental regions are exceptions characterized by relatively high AMFs. The latter are biased high from contamination of the surface albedos by persistent clouds over the monsoon regions of southeast Asia, the Ivory Coast of Africa, northern South America, New Guinea, and over the North Pacific (near 45°N, 180°E). The coefficient of variation (standard deviation divided by the monthly mean value for July 1996) of the AMF<sub>a</sub> is less than a few percent, implying that temporal variability in the shape factors from the GEOS-CHEM model has little influence on the retrieval.

[26] Figure 5 shows the relative contribution of clouds to the backscattered radiance  $I$  seen by GOME, as determined from

$$\frac{I_c f}{I} = \frac{R_c f}{R_o(1-f) + R_c f} \quad (12)$$

Generally when  $f$  is greater than 25%, over 50% of  $I$  is from the cloudy subscene. Variability in cloud optical thickness and surface albedo induces scatter in the relationship between  $I_c f / I$  and  $f$  (Figure 5). Figure 4 (middle) shows the spatial structure in  $I_c f / I$ . Clouds contribute more than 60% of  $I$  over cloudy regions such as the high-latitude oceans, the intertropical convergence zone (ITCZ), the

Asian monsoon region, and the stratus-covered oceans off the west coasts of southern Africa, South America, and North America. Clouds are less persistent over continents so that  $I_c f / I$  is relatively small and most of the backscattered radiance seen by GOME is from the clear-sky fractions of the scenes.

[27] Figure 4 (bottom) shows the actual mean AMF values for July 1996 accounting for cloud scattering as calculated with equation (11). The AMF increases by up to 40% over the oceans relative to AMF<sub>a</sub>, especially over regions where  $I$  is largely from clouds. As illustrated in Figure 2, low stratus decks over the oceans where most of the NO<sub>2</sub> is above the stratus enhance the sensitivity of GOME to the NO<sub>2</sub> column. Over continental regions with high surface emissions such as the northeastern US and northern Europe, the actual AMF is 20–30% lower than AMF<sub>a</sub> reflecting the obscuration of boundary layer NO<sub>2</sub> below clouds. The coefficient of variation of the AMF for July 1996 is generally less than 15% and always less than 30%; temporal variability in cloud cover has a moderate effect.

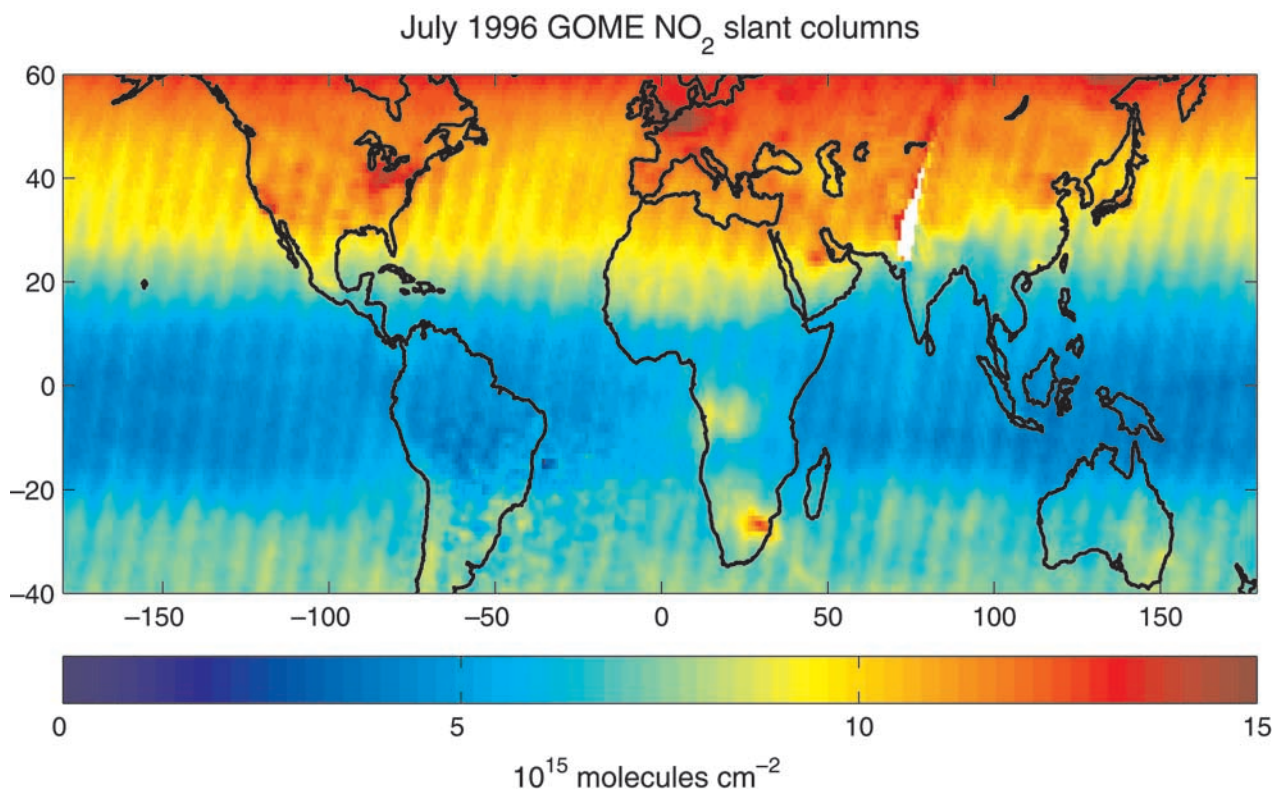
### 3.4. Comparison With Previous Retrievals

[28] The AMF calculation presented here improves in several ways over previous NO<sub>2</sub> retrievals [Leue *et al.*, 2001; Velders *et al.*, 2001; Richter and Burrows, 2002]. All three previous retrievals assumed a globally uniform NO<sub>2</sub> vertical profile in their AMF calculation. As illustrated in Figure 2, NO<sub>2</sub> vertical profiles exhibit high spatial variability, resulting in AMF values over land that are about half of ocean values. Previous retrievals also overestimated NO<sub>2</sub> columns over regions of high boundary layer depths, and underestimated them over regions of low boundary layer depths.

[29] Leue *et al.* [2001] used albedos determined over 295–745 nm, resulting in land albedos about 3 times the values used here. Both Velders *et al.* [2001] and Richter and Burrows [2002] used a globally uniform surface albedo of about 5%, closer to the values used here, but contributing to an overestimate of the NO<sub>2</sub> columns over regions of high surface albedo, and an underestimate over regions of low surface albedo. Surface albedo determined from GOME at 440 nm varies from 2% over densely vegetated regions such as the Amazon to 10–15% over deserts such as the Sahara. We find that a 50% error in the surface albedo over the Pennsylvania scene ( $0.04 \pm 0.02$ ) yields a corresponding AMF error of up to 28%. Surface albedo is particularly important in determining the sensitivity of GOME to boundary layer NO<sub>2</sub>.

[30] Richter and Burrows [2002] compared scenes with cloud fractions less than 0.1 and scenes with cloud fraction greater than 0.3 to examine the amount of NO<sub>2</sub> above clouds. Leue *et al.* [2001] and Velders *et al.* [2001] corrected for scattering by clouds by assuming a uniform cloud reflectivity of about 0.8 and cloud fraction of 0.5, a correction that increased their tropospheric values by about 4. However, cloud reflectivity is highly dependent on optical thickness and solar zenith angle (Figure 3), and is less than 0.5 for cloud optical thickness less than 10 (Figure 3, bottom right). Cloud fraction is also highly variable. Figure 4 illustrates the high spatial variability in the cloud correction from our method. We find that the actual cloud correction is much smaller than that used by Leue *et al.* [2001] and Velders *et al.* [2001] and even variable in sign, increasing tropospheric values by up to 30% over cloudy land regions but decreasing





**Figure 6.** Monthly mean slant columns of NO<sub>2</sub> determined from GOME for July 1996. Streaking along orbit paths results from daily variation in the diffuser plate artifact. This artifact is removed in section 5.

them by up to 40% over cloudy ocean regions (where most of the NO<sub>2</sub> is typically located above the cloud). Although absolute concentrations over ocean are generally small, they are important for conclusions about the amount of NO<sub>2</sub> from lightning in the free troposphere as discussed in section 6.2.

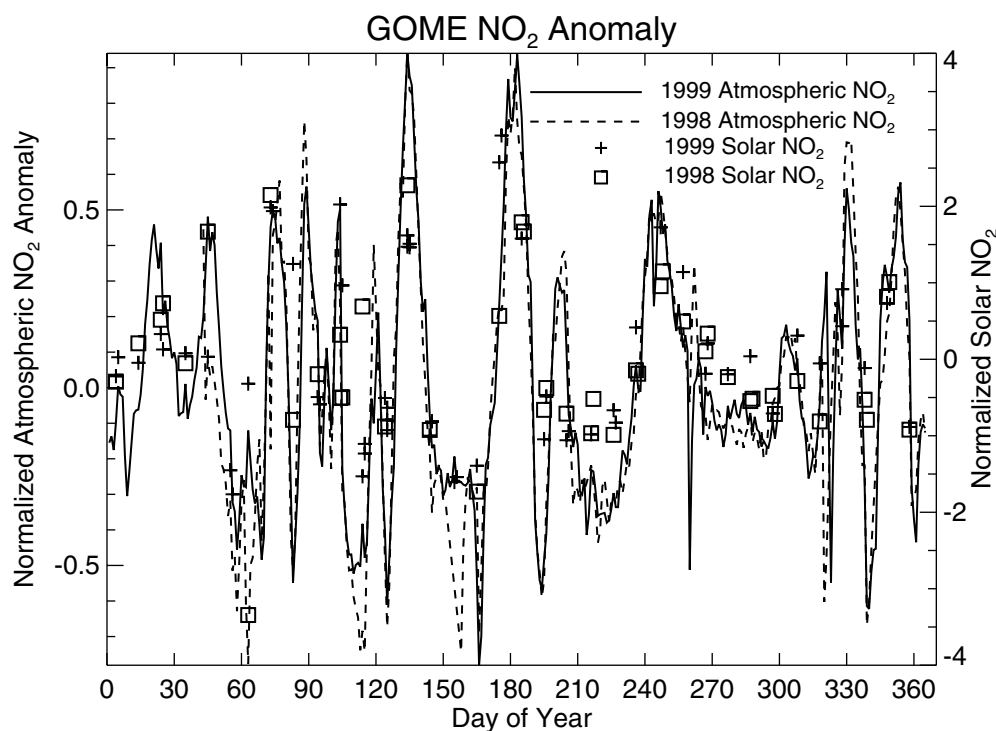
#### 4. Slant Column Fitting

[31] We determine slant columns of NO<sub>2</sub> by directly fitting backscattered radiance spectra observed by GOME, as described by *Chance* [1998]. No high-pass filtering or smoothing is applied. A nonlinear least squares inversion based on the Levenberg–Marquardt method [*Press et al.*, 1986] retrieves slant column amounts that minimize the  $\chi^2$  error between observed and calculated backscattered radiance over the wavelength region 423.08–451.23 nm. Backscattered radiances are calculated from solar irradiance spectra, measured reference spectra for the interfering species O<sub>3</sub> [*Burrows et al.*, 1999b], NO<sub>2</sub> [*Burrows et al.*, 1998], the O<sub>2</sub>–O<sub>2</sub> collision complex (P. Simon, personal communication, 1993), and H<sub>2</sub>O [*Rothman et al.*, 1998], as well as the Ring effect and the H<sub>2</sub>O-ring effect. The Ring effect is determined as described by *Chance and Spurr* [1997], including both pure Raman scattering of the Fraunhofer spectrum and the contribution from interference by atmospheric absorption. The H<sub>2</sub>O-ring effect, or the filling in of Fraunhofer lines by Raman scattering in the librational bands of liquid H<sub>2</sub>O [*Walrafen*, 1967; *Kattavar and Xu*, 1992; *Gordon*, 1999], is calculated by convolving Raman cross sections with a high resolution Fraunhofer spectrum following the work of *Chance and Spurr* [1997].

Ozone is first fitted separately between 324.93 and 335.09 nm and held at this fitted value for the NO<sub>2</sub> fitting step.

[32] The NO<sub>2</sub> spectra are available at temperatures of 221 K, 241 K, 273 K, and 293 K [*Burrows et al.*, 1998]. The magnitude of the spectral features in the 293 K NO<sub>2</sub> spectrum are about 80% of those in 221 K NO<sub>2</sub> spectrum. Slant columns fitted with the 221 K NO<sub>2</sub> spectrum are about 75–85% of those fitted with the 293 K NO<sub>2</sub> spectrum, with some spatial variability in this relationship. We use the 293 K NO<sub>2</sub> spectrum more appropriate for tropospheric NO<sub>2</sub>, which is mostly in the continental boundary layer; the resulting stratospheric NO<sub>2</sub> amounts are systematically about 20% high, but are subtracted from the total column as described below so the error is of no consequence. Previous tropospheric retrievals [*Leue et al.*, 2001; *Velders et al.*, 2001; *Richter and Burrows*, 2002] used either the 221 K or 241 K NO<sub>2</sub> spectrum more appropriate for stratospheric NO<sub>2</sub>.

[33] Wavelength calibration of the GOME backscattered radiance and solar irradiance spectra is improved using cross-correlation with the Fraunhofer reference spectrum [*Caspar and Chance*, 1997]. The aliasing introduced from severe spectral undersampling of the GOME instrument and differences between the instrument transfer functions for backscattered radiance and solar irradiance are largely corrected with an undersampling spectrum generated from the high resolution Fraunhofer reference spectrum [*Chance*, 1998]. The remaining common mode residual, determined as an average fitting residual over a complete orbit, apparently arises from instrumental artifacts (including an imperfect undersampling correction). The common mode residual



**Figure 7.** Evidence for the diffuser plate artifact. The dashed and solid lines show the normalized offset in atmospheric NO<sub>2</sub> slant columns determined from GOME for 1998 and 1999, respectively. The square and plus, are “solar NO<sub>2</sub>” fitted using the solar spectrum for 1 January 1998 as a reference. There is no known NO<sub>2</sub> in the sun.

is included in the inversion, improving the fitting precision by roughly a factor of 3. Slant columns fitted with and without the common mode residual differ by roughly  $\pm 5 \times 10^{13}$  molecules cm<sup>-2</sup>, less than 1% of the northern mid-latitude column and up to a few percent of the tropical column. The standard deviation of scene-to-scene variation in NO<sub>2</sub> slant columns over the Pacific Ocean is about half of the fitting precision determined with the common mode included and 1/6 of the fitting precision without the common mode. Accounting for the common mode residual represents better the actual fitting precision.

[34] The resulting slant columns of NO<sub>2</sub> are illustrated in Figure 6. Only observations from the forward scan are included. There are about 14 orbits daily, each at 25° intervals. Streaking along orbit paths results not from fitting errors, but rather from the observation of adjacent orbits by GOME on different days and daily variation in the diffuser plate artifact [Richter *et al.*, 2002] as discussed further below. The poleward increase arises from the stratospheric NO<sub>2</sub> column. Tropospheric signals are manifest over the northeastern US, Europe, central Africa, South Africa, and other regions. Slight speckle in the NO<sub>2</sub> columns can be seen over southern Brazil and the western South Atlantic, a region of high standard deviation in the NO<sub>2</sub> columns that is collocated with the South Atlantic Anomaly. The slant columns (determined with a 293 K NO<sub>2</sub> spectrum) are within 30% of the operational data product (determined with a 241 K NO<sub>2</sub> spectrum) [Thomas *et al.*, 1998]. The NO<sub>2</sub> fits have residual RMS values typically  $8 \times 10^{-4}$  of the GOME backscattered radiance spectrum. The mean NO<sub>2</sub> slant column fitting precision is  $1.1 \times 10^{15}$  molecules cm<sup>-2</sup>

(1σ uncertainties are used throughout). Slant columns determined with a 221 K NO<sub>2</sub> spectrum (appropriate for stratospheric NO<sub>2</sub>) have a fitting precision of  $8 \times 10^{14}$ .

[35] The diffuser plate of the GOME instrument, used to attenuate solar radiation and enable the same detector to view both the Sun and the Earth, introduces small temporally varying spectral features into the solar irradiance spectra [Richter *et al.*, 2002] (Richter and Wagner, personal communication, 2001). The spectral features vary with solar azimuth angle resulting in a temporally varying offset in the retrieved NO<sub>2</sub> slant column densities of up to  $\pm 2 \times 10^{15}$  molecules cm<sup>-2</sup>. This effect can be seen in the streaking pattern of Figure 6. Figure 7 illustrates the anomaly for 1998 and 1999 as the daily departure of retrieved NO<sub>2</sub> vertical columns from the annual mean averaged over 10°S to 10°N and 45°E to 270°E. The offset is remarkably consistent from year to year, although it has no periodicity during 1 year. Evidence that the source of the anomaly is contamination of the solar irradiance spectra by the diffuser plate can be seen in the “solar NO<sub>2</sub>” data plotted over the curves in Figure 7. These data were obtained by fitting daily GOME solar irradiance spectra for NO<sub>2</sub>, using the 1 January 1998 solar irradiance spectrum as a reference. Since NO<sub>2</sub> features are not intrinsic to actual solar irradiance spectra, we believe its apparent presence here is an artifact of the diffuser plate. Similar artifacts, not specific to GOME NO<sub>2</sub> fitting, have been observed in the laboratory and may affect all instruments that use ground aluminum diffuser plates [Richter *et al.*, 2002] (Richter and Wagner, personal communication, 2001). Other potential sources of the artifact were considered and discarded, including the GOME optical bench temper-



ature, cooler instrument warming cycles, and time-dependent level 1 calibration adjustments. Removal of the diffuser plate artifact should be performed daily, as this is the timescale for the solar irradiance measurements.

## 5. Removal of the Stratospheric Column and Diffuser Plate Artifact

[36] We define the “nontropospheric” column of NO<sub>2</sub> as the sum of the stratospheric column and instrument biases such as the diffuser plate artifact. We remove this column in a two-step process. First, we determine its amount using GOME observations over Pacific regions where tropospheric NO<sub>2</sub> is particularly low (but not zero; we subsequently correct for the resulting bias). Second, we assume it is longitudinally invariant and subtract it from the total columns in the corresponding latitude band. Errors associated with each step are discussed.

### 5.1. Determination of the Nontropospheric Column

[37] We determine the latitude-dependent nontropospheric column from GOME data over Pacific regions where tropospheric NO<sub>2</sub> is particularly low. We use the GEOS-CHEM model to identify favorable regions. Figure 8 (top) shows monthly mean tropospheric NO<sub>2</sub> slant columns calculated from the product of the AMF (Figure 4, bottom) and GEOS-CHEM tropospheric vertical columns. We determine the nontropospheric slant column for each GOME scan angle on a daily timescale and for each 2° latitude band as the zonal mean NO<sub>2</sub> column over the Pacific region within the white lines, excluding Hawaii. Scattering weights in the stratosphere are nearly independent of tropospheric clouds or surface albedo, an important consideration for being able to determine a mean stratospheric slant column; stratospheric scattering weights calculated for high reflectivity scenes ( $R = 0.8$ ) are only about 1% greater than low reflectivity scenes ( $R = 0.1$ ).

[38] The resulting monthly mean nontropospheric slant columns over the Pacific are shown in Figure 9 (top) as a function of latitude. Figure 9 (bottom) illustrates the temporal variation in the nontropospheric slant columns arising from the diffuser plate artifact for three different latitudes. There are no known photochemical or dynamical processes that could induce a factor of 2 change in equatorial NO<sub>2</sub> over a 15-day period. All latitudes generally vary together, as would be expected from an artifact that affects only the solar spectrum. The artifact decreases by up to  $5 \times 10^{14}$  molecules cm<sup>-2</sup> day<sup>-1</sup> near 7 July. Averaging the nontropospheric column for subtraction on a 2-day timescale (instead of 1 day as done here) would decrease the measurement accuracy.

[39] An alternative approach to remove the diffuser plate artifact uses a single solar spectrum in the spectral fitting [Richter and Burrows, 2002]. Like the procedure described in the current study, this method also reduces the streaking in total slant columns apparent in Figure 6. However, we find that using a single solar spectrum increases the spectral fitting uncertainty, and it is still optimal to subtract the nontropospheric column at daily intervals to reduce other instrument artifacts.

[40] We determine the nontropospheric column for each latitude band from about 17 observations over the North Pacific and up to 70 observations in the tropics. The slant

column fitting precision is typically  $1.1 \times 10^{15}$  molecules cm<sup>-2</sup> for each observation. The corresponding error decreases as  $1/\sqrt{n}$  yielding a precision between  $1.3 \times 10^{14}$  and  $2.7 \times 10^{14}$  molecules cm<sup>-2</sup>, where  $n$  is the number of observations. Determining the nontropospheric column over a narrower longitudinal or latitudinal region decreases the precision but does not significantly change the accuracy.

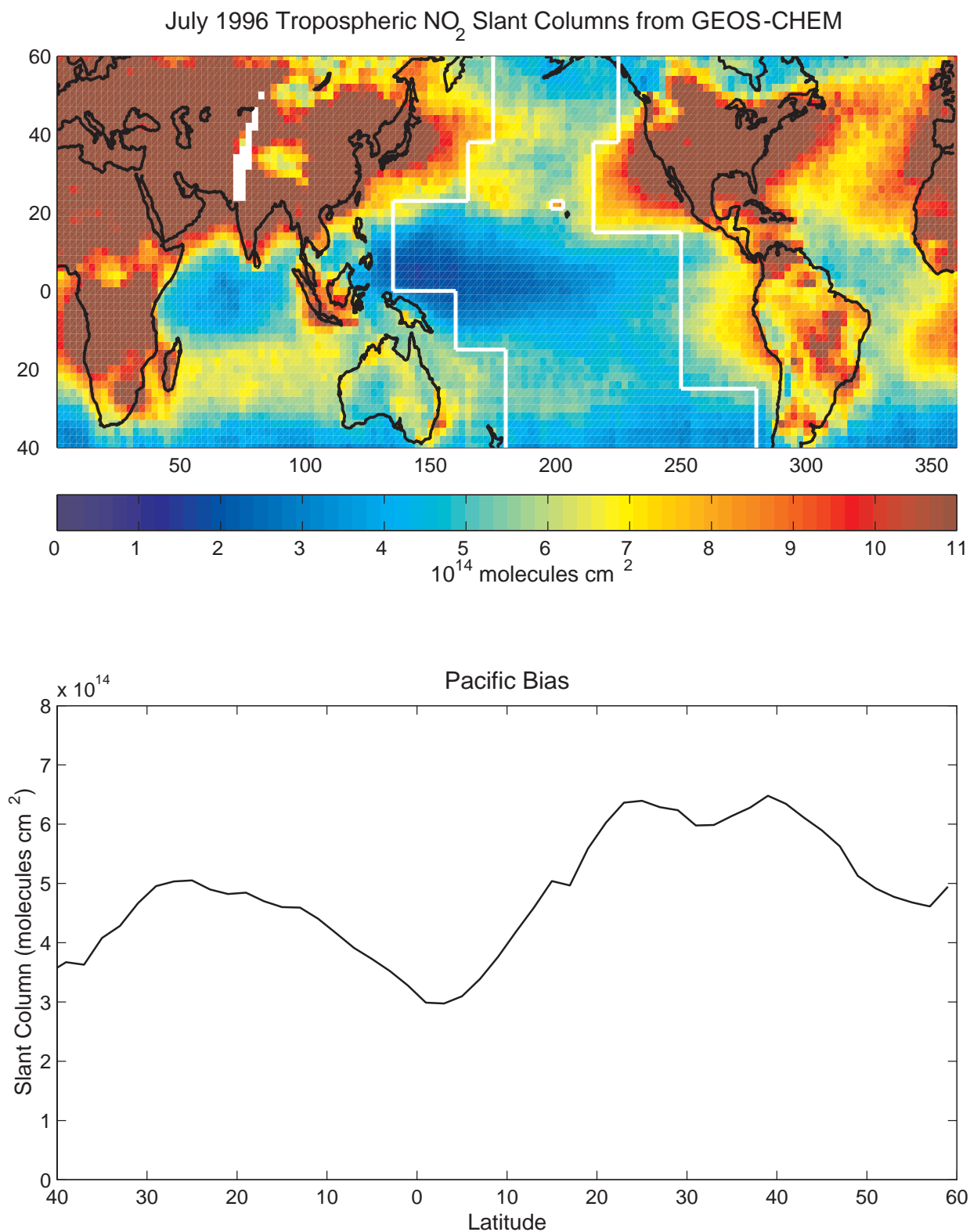
[41] Although tropospheric NO<sub>2</sub> slant columns over the Pacific are about an order of magnitude less than total slant columns, a latitudinally varying bias is introduced by the assumption of zero tropospheric NO<sub>2</sub> over the Pacific. An even larger bias would result from applying this assumption over other oceanic regions such as the North Atlantic, as done by Leue *et al.* [2001] and Velders *et al.* [2001], and is the likely cause of the negative values of tropospheric NO<sub>2</sub> columns over the oceans in their retrievals. Marine stratus decks over oceans contribute to the bias by increasing the sensitivity of GOME to tropospheric NO<sub>2</sub>.

[42] Tropospheric NO<sub>2</sub> slant columns over the Pacific (Figure 8, bottom) would introduce a latitudinally dependent bias on the nontropospheric slant column if ignored. We subtract them from the nontropospheric column derived for the corresponding latitude and for individual days using the GEOS-CHEM model results over the Pacific for that day (the monthly mean GEOS-CHEM fields are shown in the bottom panel of Figure 8). The NO<sub>2</sub> concentrations simulated by GEOS-CHEM over the Pacific are generally high (up to a factor of 2) compared with aircraft observations (e.g., Figure 1). The error introduced by the correction is discussed in section 7.

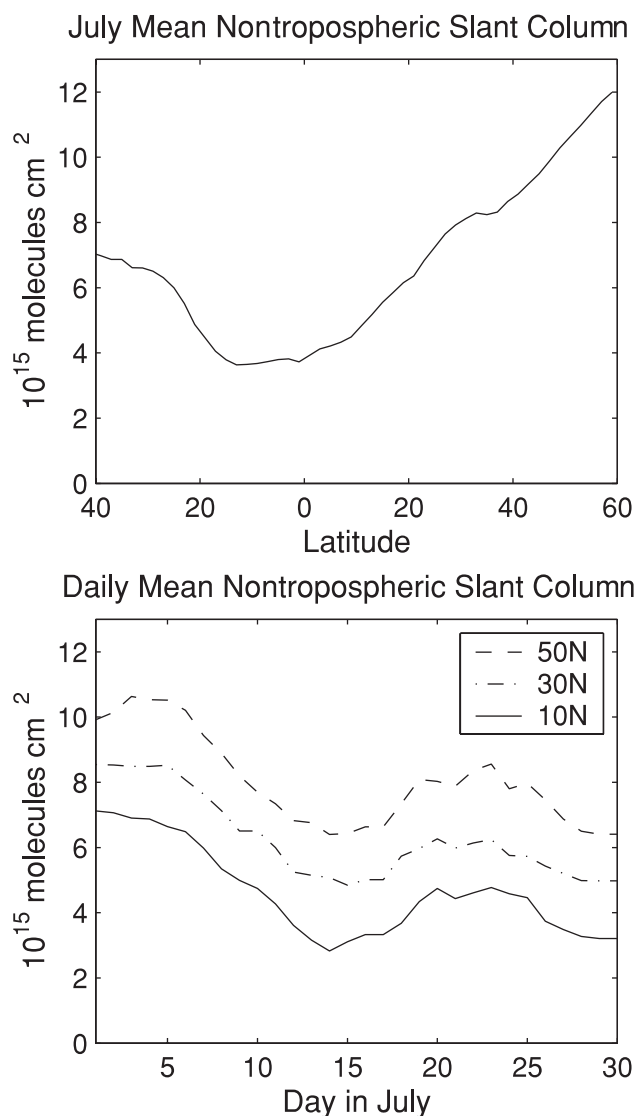
### 5.2. Zonal Invariance Assumption

[43] We subtract the nontropospheric column from each latitude for each scan angle on a daily timescale determined by each new GOME solar observation. We assume this column value to be zonally invariant. The assumption of zonal invariance introduces much less error with NO<sub>2</sub> than with ozone because most of the NO<sub>2</sub> column is in the middle stratosphere [Gordley *et al.*, 1996], whereas ozone is in the lower stratosphere and therefore more affected by zonal variability in transport and in the tropopause.

[44] We quantify the error associated with the assumption of zonally invariant stratospheric columns using NO<sub>2</sub> measurements from HALOE data [Russell *et al.*, 1993]; we use version 19 data here. Each UARS yaw cycle is about 36 days long during which time HALOE observations span approximately 60°S through 60°N. Observations by HALOE progress approximately 5° in latitude while taking about 15 observations at successive longitudes around the world. We calculate stratospheric columns of NO<sub>2</sub> from the GEOS-CHEM model tropopause to 1 hPa for each HALOE profile from yaw cycles that include the month of July for the years 1996–2000. Our analysis therefore includes observations from the end of June and the beginning of August. Years prior to 1996 are not included to reduce measurement uncertainty associated with enhanced aerosol concentrations from Mount Pinatubo [Gordley *et al.*, 1996]. Data with uncertainties greater than  $1 \times 10^8$  molecules cm<sup>-3</sup> are excluded, largely removing unreal increases near the bottom of profiles [Gordley *et al.*, 1996]. Number densities of NO<sub>2</sub> peak between 5 and 40 hPa; the resulting underestimate in stratospheric NO<sub>2</sub> columns from the



**Figure 8.** (top) Monthly mean tropospheric NO<sub>2</sub> slant columns simulated with the GEOS-CHEM model for July 1996. The white lines bound the central Pacific region of low tropospheric NO<sub>2</sub> used to determine the nontropospheric column for a given latitude. The maximum on the color scale corresponds to the GOME fitting precision and thus identifies regions where tropospheric NO<sub>2</sub> can be observed. GEOS-CHEM fields for the source regions are shown in Figure 12. (bottom) Tropospheric NO<sub>2</sub> slant columns over the central Pacific, averaged over the region bounded by the white lines and plotted as a function of latitude for July 1996. These columns represent a latitudinal bias introduced by the assumption of zero tropospheric NO<sub>2</sub> over the Pacific, and the bias is corrected as described in the text.



**Figure 9.** (top) July 1996 mean nontropospheric slant columns of NO<sub>2</sub> determined from GOME observations over the Pacific. (bottom) Daily variability in these columns determined from GOME observations over the Pacific for three different latitudes.

removal of high-uncertainty, lower-stratospheric values should be less than a few percent.

[45] We calculate means and standard deviations of the HALOE observations for each set of approximately 15 observations at different longitudes within a 5° latitude band. Figure 10 shows the standard deviations, which are a measure of zonal variability, as a function of latitude for sunrise and sunset. The standard deviation is minimum at the equator ( $1 \times 10^{14}$  molecules  $\text{cm}^{-2}$ ) and remains below  $4 \times 10^{14}$  molecules  $\text{cm}^{-2}$  at all latitudes. It is less than the GOME vertical column fitting precision, calculated as the slant column fitting precision divided by AMF<sub>G</sub>. Stratospheric columns of NO<sub>2</sub> at the time of GOME overpass (~1030 AM) are about half the sunrise/sunset values [Wennberg et al., 1994]; and the errors should be correspondingly less. The sunset values between 10°S and 25°S exhibit relatively high standard deviation across all longitudes, not

just reflecting the South Atlantic Anomaly. We find no systematic zonal pattern in the HALOE stratospheric NO<sub>2</sub> columns for latitudes equatorward of about 65°. Between 65°N and 75°N systematic maxima appear near 110°W and 80°E and minima near 20°E and 140°E, about  $\pm 4 \times 10^{14}$  molecules  $\text{cm}^{-2}$  with respect to the zonal mean (not shown).

## 6. Retrieved Tropospheric Columns of NO<sub>2</sub>

### 6.1. Slant Columns

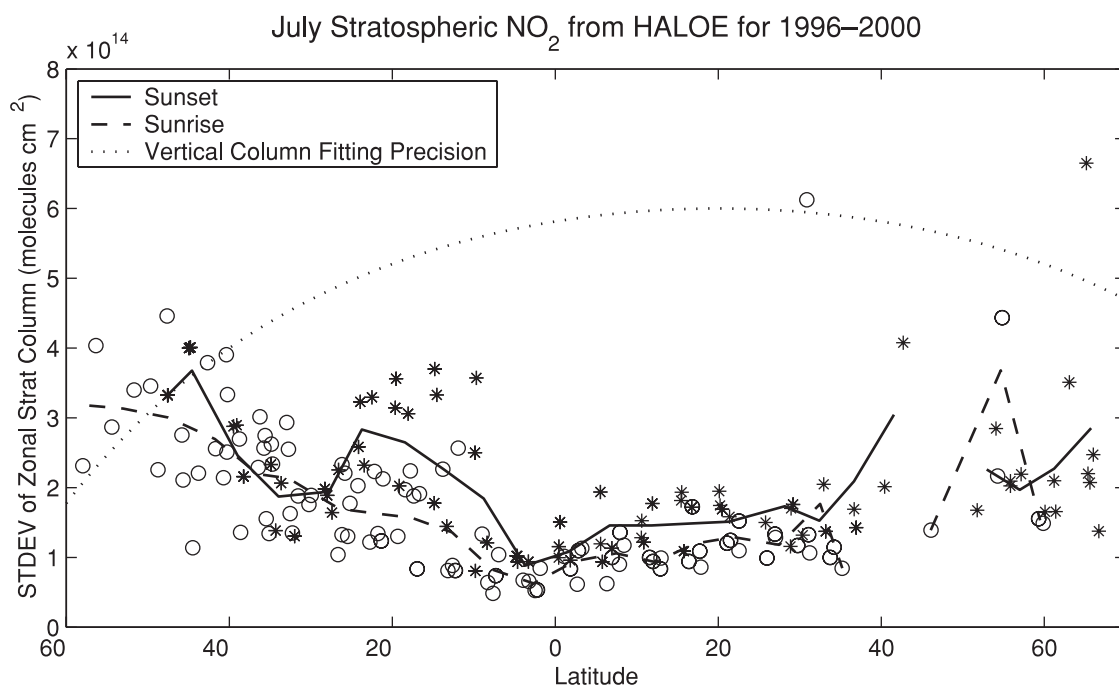
[46] Figure 11 (top) shows the monthly mean tropospheric residual slant column of NO<sub>2</sub> for July 1996. The tropospheric residual exhibits low background NO<sub>2</sub> columns over oceans and strong regional enhancements over continents. One of the strongest tropospheric enhancements is seen over the Congo Basin, where considerable biomass burning takes place in July [Scholes et al., 1996]. Over the northeastern US, California, Mexico City, and industrial regions of Europe, tropospheric enhancements are about 50% of the stratospheric column. Tropospheric enhancements are of comparable magnitude to stratospheric columns over Saudi Arabia, the biomass burning region of Central Africa, and the Transvaal region of South Africa, a major region of electricity generation. Table 2 shows that the GOME tropospheric slant columns exhibit a high degree of consistency with GEOS-CHEM vertical columns over the entire world ( $r = 0.63$ ,  $n = 7170$ ,  $p < 0.005$ ) and especially over the United States ( $r = 0.75$ ,  $n = 288$ ,  $p < 0.005$ ), even before model information has been passed to the GOME product through the AMF calculation. The Pearson's correlation coefficient is unchanged if the bias correction (Figure 8, bottom) is excluded to completely remove GEOS-CHEM influence.

[47] We have pointed out previously the sensitivity of the GOME observations to clouds, which enhance the sensitivity to NO<sub>2</sub> above clouds (as typically occurs over oceans), and obscure NO<sub>2</sub> below cloud (as typically occurs over continental source regions). The former effect is particularly obvious over the South Atlantic downwind of the Congo Basin (Figure 11), where biomass burning outflow is transported above a persistent stratus deck [Bachmeier and Fuelberg, 1996]. A similar effect over the region was noted in ozone retrievals from TOMS [Thompson et al., 1993]. Slant columns are generally higher over source regions such as the northeastern US, northern Europe, and eastern Asia where most NO<sub>2</sub> is in the boundary layer and masked by overhead clouds. We find further evidence of these cloud effects through correlations of the tropospheric NO<sub>2</sub> slant columns with GOMECAT cloud fractions. Over the northeastern United States, GOMECAT cloud fraction explains approximately 50% of the variance in tropospheric NO<sub>2</sub> slant columns ( $r = -0.71$ ,  $n = 1140$ ,  $p < 0.005$ ). In contrast, NO<sub>2</sub> columns over oceans are usually positively correlated with clouds because most NO<sub>2</sub> over oceans is in the free troposphere [Bradshaw et al., 1999] and its detection is enhanced by reflectivity from low stratus clouds. We observe this over the South Pacific, where cloud fraction explains approximately 25% of the variance in NO<sub>2</sub> tropospheric slant columns ( $r = 0.48$ ,  $n = 40377$ ,  $p < 0.005$ ).

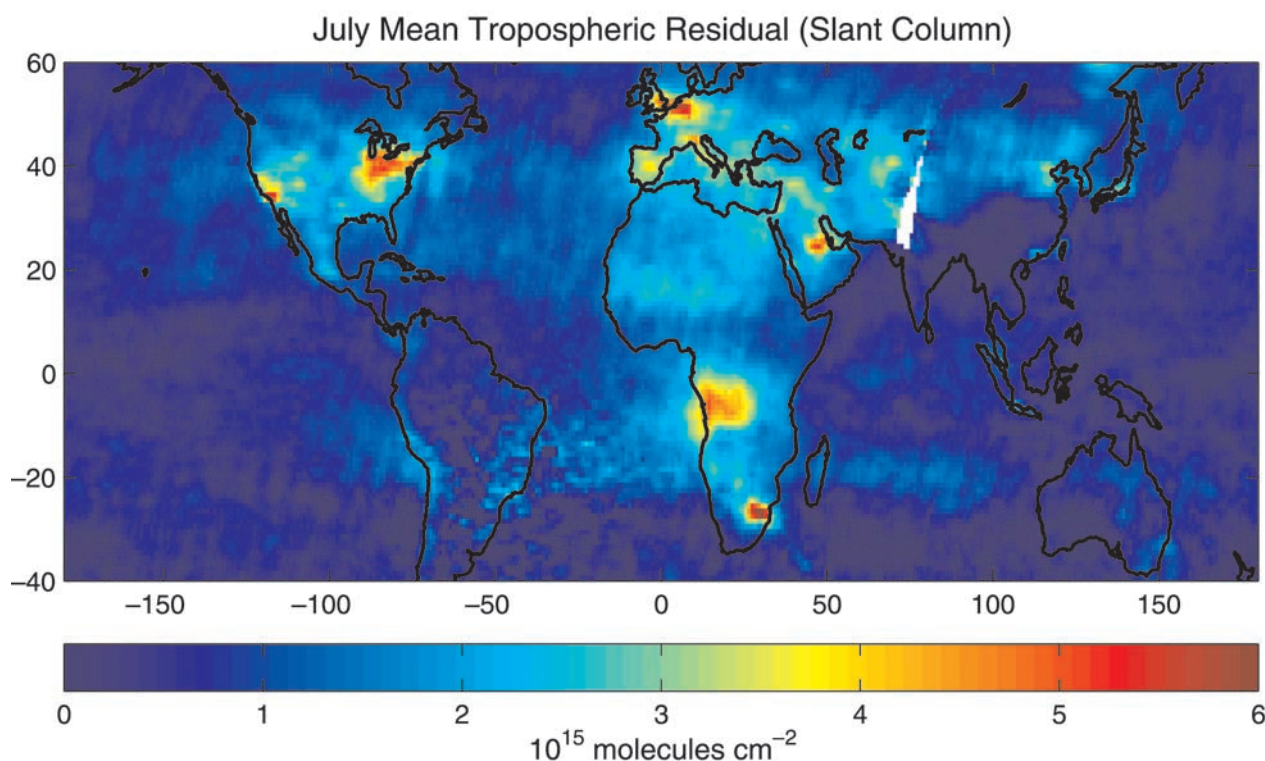
### 6.2. Vertical Columns

[48] We regrid the GOME slant columns onto the GEOS-CHEM model grid of 2° latitude by 2.5° longitude, and





**Figure 10.** Zonal standard deviation of stratospheric NO<sub>2</sub> columns (tropopause to 1 hPa) calculated from HALOE profiles for each set of roughly 15 observations in a 5° latitudinal band for each yaw cycle that includes the month of July for the years 1996–2000. Columns are calculated as described in the text. Sunset values are indicated with stars and sunrise with circles. The solid and dashed lines represent the average standard deviations. The dotted line is the vertical column fitting precision for GOME calculated as the slant column fitting precision of  $1.1 \times 10^{15}$  molecules cm<sup>-2</sup> divided by AMF<sub>G</sub>.



**Figure 11.** July 1996 mean tropospheric slant column calculated on a daily basis as the difference between total NO<sub>2</sub> slant columns and zonal mean NO<sub>2</sub> slant columns over the Pacific with bias correction as described in section 5.1.



**Table 2.** Correlation Coefficient (*r*) Between Modeled (GEOS-CHEM) and Retrieved (GOME) Tropospheric Columns<sup>a</sup>

Retrieval Step	World (n = 7170)	United States (n = 288)
Removal of the stratosphere and diffuser plate artifact	0.63	0.76
Correction for the Pacific bias	0.64	0.76
Application of AMF <sub>a</sub>	0.73	0.78
Inclusion of clouds in the AMF calculation	0.76	0.78

<sup>a</sup>All values are statistically significant (*p* < 0.005).

calculate vertical columns by dividing each tropospheric slant column by the locally derived AMF (section 3). Figure 12 (top) shows the resulting monthly mean tropospheric vertical columns. The spatial distribution is similar to that of the tropospheric slant columns (Figure 11, top). The AMF conversion from slant to vertical columns generally enhances columns over land with respect to those over oceans as discussed previously. It also enhances midlatitude columns with respect to tropical columns due to the higher solar zenith angles and shallower boundary layers at midlatitudes.

[49] Aircraft corridors or regions of intense lightning activity exhibit no NO<sub>2</sub> enhancement (Figure 12). *Richter and Burrows* [2002] attributed enhancements of NO<sub>2</sub> slant columns for scenes with cloud fractions greater than 30% over the tropical Atlantic and Africa to lightning. Although we find similar enhancements in the slant columns (Figure 11) off the west coast of the Congo Basin where persistent stratus decks make GOME particularly sensitive to NO<sub>2</sub> transported above the clouds, the enhancements largely disappear with the conversion of slant to vertical columns (Figure 12). In situ observations and the model show high NO mixing ratios in the upper troposphere over the region (top right panel of Figure 1), but the corresponding NO<sub>2</sub> number density profile is skewed toward the lower troposphere where surface NO<sub>x</sub> sources likely play a more important role. The cloud fraction threshold of 0.3 used by *Richter and Burrows* [2002] may not discriminate between upper tropospheric and surface sources since 50% of backscattered radiance still originates from the clear-sky subscene (Figure 5). To examine the sensitivity of GOME to lightning, we performed a simulation with 6 Tg NO<sub>x</sub> from lightning, twice the magnitude of the standard simulation (Table 1). The resultant GEOS-CHEM monthly mean NO<sub>2</sub> columns increased by less than  $3 \times 10^{14}$  molecules cm<sup>-2</sup> anywhere, well below the fitting uncertainty.

### 6.3. Comparison to GEOS-CHEM

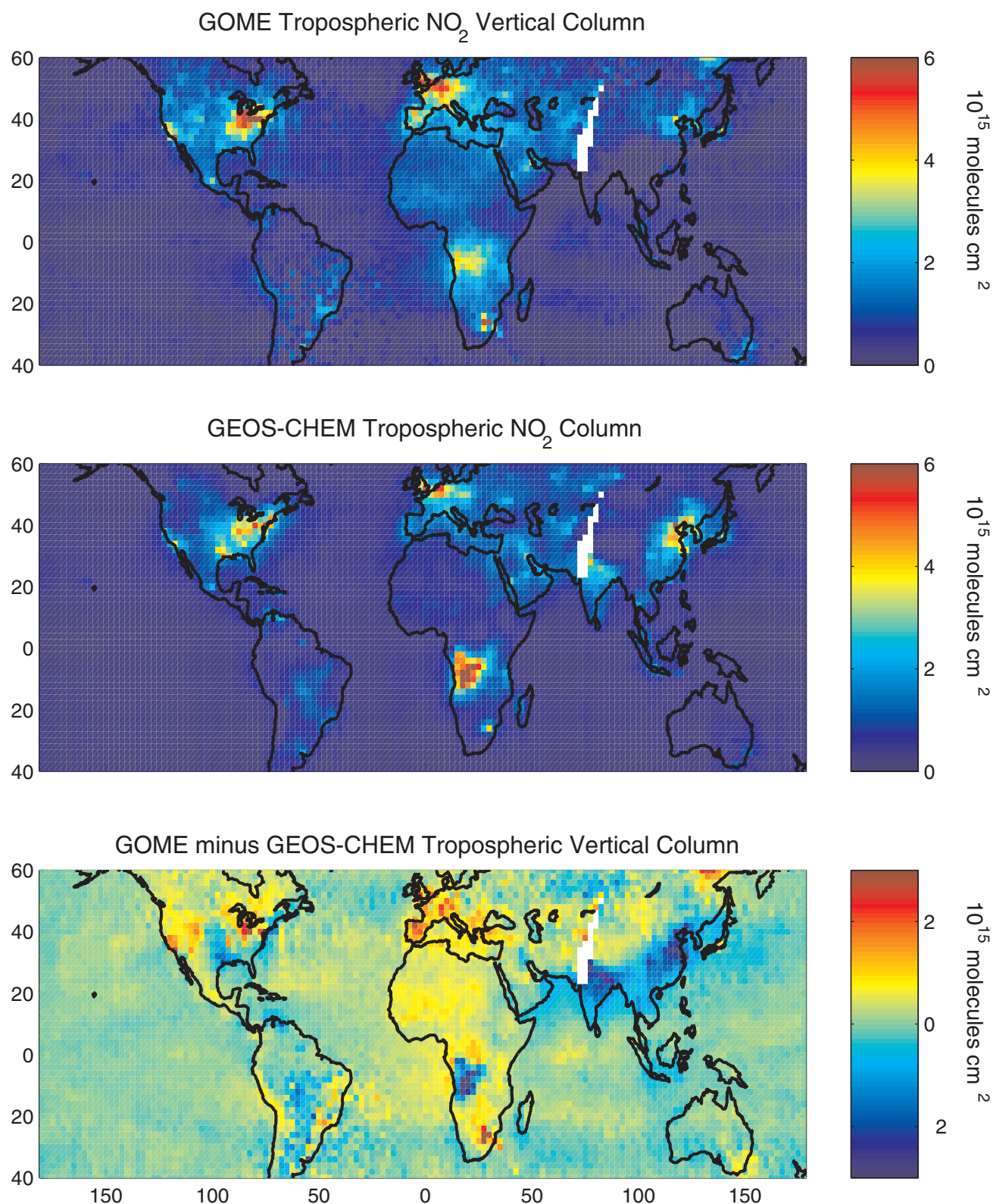
[50] A critical test of the accuracy of the GOME tropospheric NO<sub>2</sub> column is its consistency with the GEOS-CHEM model results over the United States (Figure 12, middle). The US NO<sub>x</sub> emission inventory used in GEOS-CHEM for 1996 is from the Environmental Protection Agency [EPA, 1997], and is believed to be accurate to within 20% [e.g., *Munger et al.*, 1998]. Simulation of the NO<sub>x</sub>/NO<sub>y</sub> concentration ratio at U.S. sites [*Horowitz et al.*, 1998; *Liang et al.*, 1998] suggests that the lifetime of NO<sub>x</sub> in the model is accurate to within 30%. Therefore we expect agreement between the GOME and GEOS-CHEM NO<sub>2</sub> columns to within a combined accuracy of 35%. The July mean GOME column in our retrieval over the United States

is 18% higher than the corresponding value from GEOS-CHEM, i.e., within the estimated uncertainty. The GOME vertical columns capture 58% of the spatial variance in the GEOS-CHEM columns over the United States (*r* = 0.78, *n* = 288, *p* < 0.005) as shown in Table 2. These results lend confidence to the interpretation of the GOME NO<sub>2</sub> retrieval as a proxy for surface NO<sub>x</sub> emissions. The correlation between the GEOS-CHEM model and GOME observations is also remarkable over the rest of the world (*r* = 0.76, *n* = 7170, *p* < 0.005).

[51] In comparing tropospheric NO<sub>2</sub> vertical columns from GOME with the GEOS-CHEM model, one has to be concerned about model contamination resulting from the use of GEOS-CHEM shape factors in the AMF calculation. There is really no choice in the matter, since use of shape factors from the model in the AMF calculation is a prerequisite for meaningful comparison of vertical columns [*Palmer et al.*, 2001]. Comparing slant columns would only displace the problem. One must also recognize that the shape factor and the NO<sub>2</sub> tropospheric column are two separate pieces of information, and GOME offers only one piece of information. Figure 13 shows the relationship between the AMF and modeled NO<sub>2</sub> columns, which quantifies the degree of model contamination in the comparison with observed NO<sub>2</sub> columns. Enhanced NO<sub>2</sub> columns ( $>2 \times 10^{15}$  molecules cm<sup>-2</sup>) have little relationship with the AMF (*r* = -0.14, *n* = 240) since the NO<sub>2</sub> shape factor over source regions is largely determined by boundary layer depth. Columns less than  $2 \times 10^{15}$  molecules cm<sup>-2</sup> exhibit a significant relationship with the AMF (*r* = -0.65, *n* = 6987, *p* < 0.005) that arises in part from the contrast between NO<sub>2</sub> vertical profiles over ocean and land.

[52] Retrieved tropospheric vertical columns exhibit a number of differences with the GEOS-CHEM model columns (Figure 12, bottom). Observations from GOME indicate about 50% more NO<sub>2</sub> over the Transvaal region of South Africa, a major region of electricity generation. Retrieved columns also are higher over the northeastern US and industrial regions of Europe. Observations show less NO<sub>2</sub> over Houston, the biomass burning region of central Africa, northern India, and eastern Asia. Neglect of absorbing aerosols in the AMF calculation could contribute to a retrieval underestimate for the latter regions, where soot concentrations are particularly high. A similar problem was reported by *Fishman et al.* [1996a] in retrieving tropospheric ozone columns from TOMS over biomass burning regions.

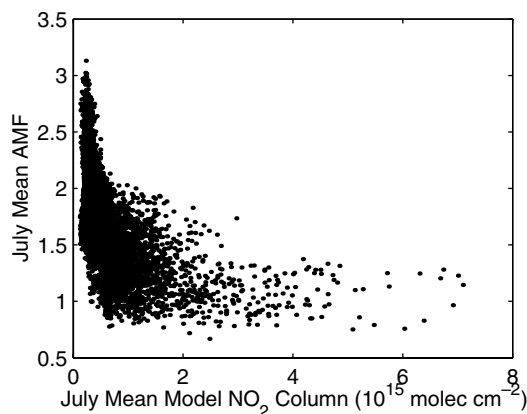
[53] The spatial distribution of GOME tropospheric NO<sub>2</sub> vertical columns presented by *Velders et al.* [2001] for July 1997 exhibits similar enhancements over industrial regions



**Figure 12.** July 1996 mean tropospheric NO<sub>2</sub> vertical columns from GOME (top) and the GEOS-CHEM model (middle). (bottom) The difference between the two (note change of scale).

of Europe, the eastern United States, eastern Asia, and over the biomass burning region of central Africa, but their NO<sub>2</sub> columns over each region are about twice the magnitude of the NO<sub>2</sub> columns presented here for July 1996. As dis-

cussed in section 3.4, this difference can be explained by their treatment of clouds in the AMF calculation. Another difference is their assumption of zero tropospheric NO<sub>2</sub> over oceanic scenes, which we argue leads to significant bias



**Figure 13.** Relationship between the AMF and the modeled NO<sub>2</sub> columns for July 1996.

(section 5.1). Comparison between the NO<sub>2</sub> columns presented here with those presented by *Leue et al.* [2001] is more difficult since they only provide yearly mean values and tropospheric NO<sub>2</sub> columns in winter are about twice the magnitude of those in summer [*Velders et al.*, 2001]. The spatial structure between the two exhibit some consistency over land.

## 7. Error Analysis

[54] In the previous sections we assessed the errors introduced at different steps of the retrieval. Here we synthesize this information to estimate the total error and identify the dominant contributions. We express the total tropospheric vertical column error  $\epsilon$  as the quadrature sum of (1) the slant column fitting error  $\epsilon_f$ , (2) the error  $\epsilon_s$  incurred in the determination of the nontropospheric slant column, (3) the error  $\epsilon_b$  in the bias correction for tropospheric NO<sub>2</sub> over the Pacific, (4) the error  $\epsilon_v$  from zonal variability in the stratospheric column, (5) the error  $\epsilon_\sigma$  due to uncertainty in the NO<sub>2</sub> absorption cross section, and (6) the error  $\epsilon_A$  in the AMF calculation

$$\epsilon = \sqrt{\left(\frac{\epsilon_f}{\text{AMF}}\right)^2 + \left(\frac{\epsilon_s}{\text{AMF}}\right)^2 + \left(\frac{\epsilon_b}{\text{AMF}}\right)^2 + \epsilon_v^2 + (\epsilon_\sigma \Omega)^2 + (\epsilon_A \Omega)^2}. \quad (13)$$

The slant column errors in the first three terms are normalized by the tropospheric AMF and the last two are normalized by the tropospheric vertical column  $\Omega$ .

[55] Table 3 summarizes the error estimates. The tropospheric AMF is taken here as 1.3 over land and 2.0 over ocean (Figure 4, bottom). The slant column fitting precision  $\epsilon_f$  is typically  $1.1 \times 10^{15}$  molecules cm<sup>-2</sup>. Over northern midlatitudes where the minimum number of observations (17) is used to determine the nontropospheric slant columns, the error  $\epsilon_s$  is  $3 \times 10^{14}$  molecules cm<sup>-2</sup>. The GEOS-CHEM simulation of aircraft observations over the remote Pacific can overestimate NO<sub>2</sub> columns by up to a factor of 2; at 40°N where the bias correction for tropospheric NO<sub>2</sub> is largest (Figure 8, bottom), the slant column error  $\epsilon_b$  is about  $3.5 \times 10^{14}$  molecules cm<sup>-2</sup>. Figure 10 shows that poleward of 30°  $\epsilon_v$  is about  $3 \times 10^{14}$  molecules cm<sup>-2</sup> at sunrise and sunset. At the time of the GOME overpass (~1030 AM),  $\epsilon_v$  is about half the sunrise/sunset values, small compared with  $\epsilon_f$ . The error in the NO<sub>2</sub> cross section is 4% [*Burrows et al.*, 1998], negligible over both land and ocean.

[56] We calculate  $\epsilon_a$  separately over ocean and land. Over ocean the tropospheric NO<sub>2</sub> column is typically less than  $2 \times 10^{15}$  molecules cm<sup>-2</sup>. Clouds are the most important contributor to variability in the AMF there, with typically a 30% effect. Even if a 50% error exists in GOMECAT cloud data,  $\epsilon_A \Omega$  would be  $3 \times 10^{14}$  molecules cm<sup>-2</sup>, small compared to the error from  $\epsilon_f$ .

[57] Over land the tropospheric NO<sub>2</sub> column ranges from  $1 \times 10^{15}$  to  $6 \times 10^{15}$  molecules cm<sup>-2</sup> and the relative error in the AMF is more important. The AMF calculation is most sensitive to surface albedo, the NO<sub>2</sub> shape factor, and cloud information. The precision of the surface albedo data is 0.02 [*Koelemeijer et al.*, 2002], yielding a corresponding error  $\epsilon_a$  of up to 28% over the Pennsylvania scene for a typical surface albedo of 0.04. Surface albedo is particularly important in determining the sensitivity of GOME to boundary layer NO<sub>2</sub> over non-frozen surfaces.

[58] Most NO<sub>2</sub> over land is in the boundary layer and the shape of the NO<sub>2</sub> profile is largely determined by boundary layer depth. As shown by *Fiore et al.* [2002], GEOS July mean boundary layer depths used in GEOS-CHEM are consistent with observations over the United States. Temporal variability is more difficult to evaluate given the limited number of observations. Inland observations over the northeastern United States indicate that daily variability in the July boundary layer depth is less than 100 hPa [*Berman et al.*, 1999]. We calculate the error in the AMF calculation by assuming that the GEOS data have no skill in simulating  $\pm 100$  hPa. For the Pennsylvania case (Figure 2), we find that AMFs calculated for boundary

**Table 3.** Error in the Retrieval of Tropospheric Vertical NO<sub>2</sub> Column From GOME (10<sup>15</sup> molecules cm<sup>-2</sup>)

	Ocean	Land
Fitting, $\epsilon_f/\text{AMF}$	0.6	0.8
Determination of the nontropospheric column, $\epsilon_s/\text{AMF}$	0.1	0.2
Pacific bias correction, $\epsilon_b/\text{AMF}$	0.2	0.3
Assumption of a zonally invariant stratospheric column, $\epsilon_v$	0.2	0.2
NO <sub>2</sub> cross section, $\epsilon_\sigma \Omega$	0.04	0.04–0.2
AMF calculation, $\epsilon_A \Omega$	0.3	0.5–3.2
<b>Total error (quadrature sum)</b>	<b>0.7</b>	<b>1.0–3.3</b>

The above errors are for each observation. In the unlikely event of truly random errors, the monthly mean error for each  $2^\circ \times 2.5^\circ$  model grid box would be  $1 \times 10^{14}$  molecules cm<sup>-2</sup> over ocean and  $2\text{--}7 \times 10^{14}$  molecules cm<sup>-2</sup> over land.



layer depths of 700 and 900 hPa differ by 15% from the value at 800 hPa.

[59] Clouds and aerosols are an obvious source of error in the AMF calculation. Of the three cloud parameters used to calculate AMFs in the cloudy scenes (cloud fraction  $f$ , cloud top pressure, and cloud optical thickness  $\tau_c$ ) the cloud optical thickness makes the largest contribution to the AMF error. We find that cloud top pressure generally has little effect on the AMF over land since they are nearly always above the boundary layer. Both  $f$  and  $\tau_c$  affect the fraction of  $I$  from the cloudy subscene in the AMF calculation (equation (11)), but of the two  $\tau_c$  is more uncertain. The error in the  $\tau_c$  measurement increases with decreasing  $f$ . We make a conservative estimate of the error in the AMF calculation by assuming factors of 10, 5, and 2 errors in  $\tau_c$  for the Pennsylvania scene for  $f$  of 0.1, 0.2, and 0.5, respectively, yielding corresponding errors of 16%, 28%, and 27%. The error in the AMF decreases with decreasing  $f$  in spite of the increased error in  $\tau_c$ . In fact the errors on  $f$  and  $\tau_c$  derived by GOMECAT are negatively correlated because of the constraint from observed reflectivity; therefore the combined error in the AMF from  $f$  and  $\tau_c$  is generally less than from  $\tau_c$  alone. We estimate an additional 10% error from uncertainties in modeling the cloud radiation transfer characteristics such as the phase function and the representation of multiple clouds as a single cloud. As pointed out earlier, the GOMECAT algorithm effectively treats a thick aerosol layer as a thin cloud. Palmer *et al.* [2001] found a 30% decrease of the AMF for a scattering aerosol in the boundary layer with an optical thickness of 1 [Palmer *et al.*, 2001]. Absorbing aerosols would have an opposite effect. Combining all the above effects, an upper bound for the quadrature sum of the error from surface albedo, boundary layer depths, clouds, and aerosols is 53%. The corresponding total error  $\epsilon_A \Omega$  ranges from  $5 \times 10^{14}$  to  $3.2 \times 10^{15}$  molecules cm<sup>-2</sup> for a tropospheric NO<sub>2</sub> column range of  $1 \times 10^{15}$  to  $6 \times 10^{15}$  molecules cm<sup>-2</sup>.

[60] As summarized in Table 3, the total error on the NO<sub>2</sub> tropospheric column retrieval is dominated by the fitting precision over ocean and land regions with low NO<sub>2</sub> columns. Over major continental source regions, the AMF calculation can be a more important contributor to the total error, due mostly to errors in surface albedo and clouds. On average about 2.5 observations are made in each  $2^\circ \times 2.5^\circ$  grid box for each of the 10 GOME overpasses per month. If the errors were truly random, the monthly mean error for each model grid box would be  $1 \times 10^{14}$  molecules cm<sup>-2</sup> over ocean and  $2\text{--}7 \times 10^{14}$  molecules cm<sup>-2</sup> over land. In practice, the errors likely include some systematic biases such as from surface albedo and aerosols.

## 8. Conclusions and Recommendations for Further Work

[61] We have presented a retrieval of tropospheric NO<sub>2</sub> vertical columns from GOME that improves in several ways over previous retrievals, especially in the AMF formulation used to convert slant columns to vertical columns. For each GOME observation, we calculate an AMF from the relative

vertical NO<sub>2</sub> distribution (shape factor) determined locally with a 3-D global model of tropospheric chemistry (GEOS-CHEM), weighted by altitude-dependent scattering weights computed with a radiative transfer model (LIDORT). The AMF calculation uses local surface albedos determined from GOME near the center of the NO<sub>2</sub> fitting window (440 nm) and near the same wavelength used in the AMF calculation. It accounts for cloud scattering using GOME cloud fraction, cloud top pressure, and cloud optical thickness from a cloud retrieval algorithm (GOMECAT). We find that clouds increase the sensitivity of GOME to tropospheric NO<sub>2</sub> columns over ocean by up to 40%, and decrease the sensitivity of GOME to tropospheric NO<sub>2</sub> columns over continental source regions by 20–30%. In general GOME is almost twice as sensitive to tropospheric NO<sub>2</sub> columns over ocean than over land due to differences in the shape of the NO<sub>2</sub> profile.

[62] Several additional algorithm improvements were presented. Slant NO<sub>2</sub> columns, which are directly fitted without low-pass filtering or spectral smoothing, are corrected daily for an artificial offset likely induced by spectral structure on the diffuser plate of the GOME instrument. We determine the stratospheric NO<sub>2</sub> column and the magnitude of the diffuser plate artifact by using data over the central Pacific Ocean, where tropospheric NO<sub>2</sub> is particularly low, and then use the GEOS-CHEM model to account for non-zero tropospheric NO<sub>2</sub> over that region. Retrieved columns are available at <http://www-as.harvard.edu/chemistry/trop/satellite/no2.html>.

[63] Retrieved tropospheric vertical columns from GOME for July 1996 exhibit a high degree of consistency with simulated columns from GEOS-CHEM. Over the United States, where NO<sub>x</sub> emissions are particularly well known, retrieved columns are 18% higher than GEOS-CHEM and monthly mean values are strongly spatially correlated ( $r = 0.78$ ,  $n = 288$ ,  $p < 0.005$ ). The GOME columns are lower than GEOS-CHEM columns over Houston, the biomass burning region of central Africa, northern India, and eastern Asia; all those have high concentrations of absorbing aerosols that could introduce a bias in the AMF calculation. Retrieved columns over the Transvaal region of South Africa, a major region of electricity generation, are about 50% higher than GEOS-CHEM values. GOME columns also are higher over the northeastern US and industrial regions of Europe. No enhancements are apparent over aircraft corridors or regions of intense lightning activity, consistent with the strong weighting of the NO<sub>2</sub> vertical distribution toward the lower troposphere.

[64] This study confirms that GOME observations can be used to map surface emissions of NO<sub>2</sub>. Lightning activity is much more difficult to detect due to large relative errors over oceans, the weak sensitivity of NO<sub>2</sub> columns to upper tropospheric NO<sub>x</sub>, and the large GOME spatial resolution. The error in the retrieval of tropospheric NO<sub>2</sub> is dominated by the spectral fitting precision over the oceans and over continental regions with low NO<sub>2</sub> columns. Over regions of enhanced NO<sub>2</sub> columns ( $>2 \times 10^{15}$  molecules cm<sup>-2</sup>) the AMF calculation becomes a more important contributor to the total error mostly because of clouds, aerosols, and surface albedo.

[65] Extension of this analysis to a full year should examine seasonal changes in zonal variability of strato-



spheric NO<sub>2</sub> columns and daily varying surface albedos from snow cover. Accounting for aerosols should improve the accuracy of the AMF calculation since NO<sub>2</sub> emissions are frequently associated with aerosols. Independent characterization of temporal variation in the diffuser plate artifact may enable more accurate removal of the stratospheric vertical columns. Future satellite instruments with smaller fields of view to reduce cloud contamination, such as SCIAMACHY on board ENVISAT and OMI on board AURA, should further improve the potential to retrieve tropospheric NO<sub>2</sub> columns from solar backscatter measurements.

[66] **Acknowledgments.** We gratefully acknowledge helpful discussions about HALOE and stratospheric variability with Anne Douglass, Dylan Jones, Ross Salawitch, and Hans Schneider. We thank Folkert Boersma, Ellen Brinksma, John Burrows, Ulrich Platt, Andreas Richter, Piet Stammes, and Pepijn Veekind for comments. This work was supported by the NASA Atmospheric Chemistry Modeling and Analysis Program. Randall V. Martin was supported by a National Defense and Engineering Graduate Fellowship.

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